



Approximation for Inextensible Flows of Curves in \mathbb{E}^3

Talat Körpınar and Essin Turhan

ABSTRACT: In this paper, we construct a new method for inextensible flows of curves in \mathbb{E}^3 . Using the Frenet frame of the given curve, we present partial differential equations. We give some characterizations for curvatures of a curve in \mathbb{E}^3 .

Key Words: Fluid flow, \mathbb{E}^3 , partial differential equation.

Contents

1 Introduction	45
2 Preliminaries	46
3 Inextensible Flows of Curves in \mathbb{E}^3	47

1. Introduction

Flows has received considerable recent attention because of its relatively simple geometry and its relevance to a large variety of engineering applications [20]. It enables direct insight into fundamental turbulence physics, as well as direct verification by local (such as profiles of mean velocity, Reynolds stress, fluctuation intensities, etc.) and integral measurements (such as skin friction and heat transfer). The fundamental challenge is to predict the mean flow properties, including the mean velocity profile (MVP), the Reynolds stress, the kinetic energy, etc; however, a deductive theory of this kind is still missing, [17].

In the past two decades, for the need to explain certain physical phenomena and to solve practical problems, geometers and geometric analysis have begun to deal with curves and surfaces which are subject to various forces and which flow or evolve with time in response to those forces so that the metrics are changing. Now, various geometric flows have become one of the central topics in geometric analysis. Many authors have studied geometric flow problems, [12].

This study is organised as follows: Firstly, we study inextensible flows of curves in Euclidean 3-space. Secondly, using the Frenet frame of the given curve, we present partial differential equations. Finally, we give some characterizations for curvatures of a curve in Euclidean 3-space.

2. Preliminaries

To meet the requirements in the next sections, here, the basic elements of the theory of curves in the space \mathbb{E}^3 are briefly presented; a more complete elementary treatment can be found in [8].

The Euclidean 3-space \mathbb{E}^3 provided with the standard flat metric given by

$$\langle, \rangle = dx_1^2 + dx_2^2 + dx_3^2,$$

where (x_1, x_2, x_3) is a rectangular coordinate system of \mathbb{E}^3 . Recall that, the norm of an arbitrary vector $a \in \mathbb{E}^3$ is given by

$$\|a\| = \sqrt{\langle a, a \rangle}.$$

α is called a unit speed curve if velocity vector v of α satisfies $\|v\| = 1$. Let $\alpha = \alpha(s)$ be a regular curve in \mathbb{E}^3 . If the tangent vector of this curve forms a constant angle with a fixed constant vector U , then this curve is called a general helix or an inclined curve. The sphere of radius $r > 0$ and with center in the origin in the space \mathbb{E}^3 is defined by

$$\mathbb{S}^2 = \{p = (p_1, p_2, p_3) \in \mathbb{E}^3 : \langle p, p \rangle = r^2\}.$$

Denote by $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$ the moving Frenet-Serret frame along the curve α in the space \mathbb{E}^3 . For an arbitrary curve α with first and second curvature, κ and τ in the space \mathbb{E}^3 , the following Frenet-Serret formulae are given in [8] written under matrix form

$$\begin{bmatrix} \mathbf{T}' \\ \mathbf{N}' \\ \mathbf{B}' \end{bmatrix} = \begin{bmatrix} 0 & \kappa & 0 \\ -\kappa & 0 & \tau \\ 0 & -\tau & 0 \end{bmatrix} \begin{bmatrix} \mathbf{T} \\ \mathbf{N} \\ \mathbf{B} \end{bmatrix},$$

where

$$\begin{aligned} \langle \mathbf{T}, \mathbf{T} \rangle &= \langle \mathbf{N}, \mathbf{N} \rangle = \langle \mathbf{B}, \mathbf{B} \rangle = 1, \\ \langle \mathbf{T}, \mathbf{N} \rangle &= \langle \mathbf{N}, \mathbf{B} \rangle = \langle \mathbf{T}, \mathbf{B} \rangle = 0. \end{aligned}$$

Let $\mathbf{u} = (u_1, u_2, u_3)$, $\mathbf{v} = (v_1, v_2, v_3)$ and $\mathbf{w} = (w_1, w_2, w_3)$ be vectors in \mathbb{E}^3 and $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$ be positive oriented natural basis of \mathbb{E}^3 . Cross product of \mathbf{u} and \mathbf{v} is defined by

$$\mathbf{u} \times \mathbf{v} = \begin{vmatrix} \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \\ u_1 & u_2 & u_3 \\ v_1 & v_2 & v_3 \end{vmatrix}.$$

Mixed product of \mathbf{u} , \mathbf{v} and \mathbf{w} is defined by the determinant

$$[\mathbf{u}, \mathbf{v}, \mathbf{w}] = \begin{vmatrix} u_1 & u_2 & u_3 \\ v_1 & v_2 & v_3 \\ w_1 & w_2 & w_3 \end{vmatrix}.$$

Torsion of the curve α is given by the aid of the mixed product

$$\tau = \frac{[\alpha', \alpha'', \alpha''']}{\kappa^2}.$$

3. Inextensible Flows of Curves in \mathbb{E}^3

Physically, inextensible curve and surface flows give rise to motions in which no strain energy is induced. The swinging motion of a cord of fixed length, for example, or of a piece of paper carried by the wind, can be described by inextensible curve and surface flows. Such motions arise quite naturally in a wide range of physical applications, [11,12,13].

Let $\alpha(u, t)$ is a one parameter family of smooth curves in \mathbb{E}^3 .

Any flow of α can be represented as

$$\frac{\partial \alpha}{\partial t} = f_1 \mathbf{T} + f_2 \mathbf{N} + f_3 \mathbf{B}, \quad (3.1)$$

where f_1, f_2, f_3 are smooth functions.

Letting the arclength variation be

$$s(u, t) = \int_0^u v du.$$

In the \mathbb{E}^3 the requirement that the curve not be subject to any elongation or compression can be expressed by the condition

$$\frac{\partial}{\partial t} s(u, t) = \int_0^u \frac{\partial v}{\partial t} du = 0.$$

Definition 3.1. The flow $\frac{\partial \alpha}{\partial t}$ in \mathbb{E}^3 are said to be inextensible if

$$\frac{\partial}{\partial t} \left\| \frac{\partial \alpha}{\partial u} \right\| = 0. \quad (3.2)$$

Theorem 3.2. Let $\frac{\partial \alpha}{\partial t} = f_1 \mathbf{T} + f_2 \mathbf{N} + f_3 \mathbf{B}$ be a smooth flow of α . The flow is inextensible if and only if

$$\frac{\partial f_1}{\partial u} = \kappa f_2 v. \quad (3.3)$$

Proof: Assume that $\frac{\partial \alpha}{\partial u}$ is inextensible. Then,

$$\frac{\partial}{\partial t} \sigma(u, t) = \int_0^u \frac{\partial v}{\partial t} du = \int_0^u \left[\frac{\partial f_1}{\partial u} - \kappa f_2 v \right] du = 0. \quad (3.4)$$

Substituting (3.2) in (3.4) complete the proof of the theorem. \square

We now restrict ourselves to arc length parametrized curves. That is, $v = 1$ and the local coordinate u corresponds to the curve arc length s . We require the following lemma.

Lemma 3.3.

$$\frac{\partial \mathbf{T}}{\partial t} = [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] \mathbf{N} + [\frac{\partial f_3}{\partial s} + f_2 \tau] \mathbf{B}, \quad (3.5)$$

where f_1, f_2, f_3 are smooth functions of time and arc length.

Proof: Using definition of α , we have

$$\frac{\partial}{\partial t} \mathbf{T} = \left[\frac{\partial f_1}{\partial s} - \kappa f_2 \right] \mathbf{T} + [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] \mathbf{N} + [\frac{\partial f_3}{\partial s} + f_2 \tau] \mathbf{B}. \quad (3.6)$$

Substituting (3.3) in (3.6), we obtain (3.5). This completes the proof. \square

Now we give the characterization of evolution of curvature as below:

Theorem 3.4. Let $\frac{\partial \alpha}{\partial t}$ be inextensible. Then, the evolution of curvature

$$\frac{\partial \kappa}{\partial t} = \frac{\partial}{\partial s} [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] - \tau [\frac{\partial f_3}{\partial s} + f_2 \tau],$$

where f_1, f_2, f_3 are smooth functions of time and arc length.

Proof: Assume that $\frac{\partial \alpha}{\partial t}$ is inextensible in \mathbb{E}^3 .

Thus it is easy to obtain that

$$\begin{aligned} \frac{\partial}{\partial s} \frac{\partial}{\partial t} \mathbf{T} &= -\kappa [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] \mathbf{T} \\ &\quad + \left[\frac{\partial}{\partial s} [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] - \tau [\frac{\partial f_3}{\partial s} + f_2 \tau] \right] \mathbf{N} \\ &\quad + \left[\frac{\partial}{\partial s} [\frac{\partial f_3}{\partial s} + f_2 \tau] + \tau [f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau] \right] \mathbf{B}. \end{aligned} \quad (3.7)$$

On the other hand, we have

$$\frac{\partial \kappa}{\partial t} = \frac{\partial}{\partial t} \langle \frac{\partial}{\partial s} \mathbf{T}, \mathbf{N} \rangle.$$

Also,

$$\frac{\partial \kappa}{\partial t} = \langle \frac{\partial}{\partial t} \frac{\partial}{\partial s} \mathbf{T}, \mathbf{N} \rangle + \langle \frac{\partial}{\partial s} \mathbf{T}, \frac{\partial}{\partial t} \mathbf{N} \rangle.$$

From definition of flow, we have

$$\langle \mathbf{N}, \frac{\partial}{\partial t} \mathbf{N} \rangle = 0.$$

Combining these we have

$$\frac{\partial \kappa}{\partial t} = \frac{\partial}{\partial s} [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] - \tau [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau].$$

Thus, we obtain the theorem. This completes the proof. \square

From the above theorem, we have

Theorem 3.5.

$$\frac{\partial \mathbf{N}}{\partial t} = -[\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \mathbf{T} + \frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]] \mathbf{B}.$$

Proof: Using Frenet equations, we have

$$\frac{\partial}{\partial t} \frac{\partial}{\partial s} \mathbf{T} = \frac{\partial}{\partial t} (\kappa \mathbf{N}) = \frac{\partial \kappa}{\partial t} \mathbf{N} + \kappa \frac{\partial \mathbf{N}}{\partial t}. \quad (3.8)$$

Substituting (3.7) into Eq. (3.8), we have

$$\begin{aligned} \frac{\partial \mathbf{N}}{\partial t} &= -[\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \mathbf{T} \\ &\quad + \frac{1}{\kappa} [\frac{\partial}{\partial s} [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] - \tau [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] - \frac{\partial \kappa}{\partial t}] \mathbf{N} \\ &\quad + \frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]] \mathbf{B}. \end{aligned}$$

On the other hand, using above equation we have

$$\frac{\partial \mathbf{N}}{\partial t} = -[\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \mathbf{T} + \frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]] \mathbf{B},$$

which completes the proof. \square

Now we give the characterization of evolution of torsion as below:

Theorem 3.6. Let $\frac{\partial \alpha}{\partial t}$ be inextensible. Then, the evolution of curvature

$$\frac{\partial \tau}{\partial t} = \frac{\partial}{\partial s} [\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]] + \kappa [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau],$$

where $\mathfrak{f}_1, \mathfrak{f}_2, \mathfrak{f}_3$ are smooth functions of time and arc length.

Proof:

$$\begin{aligned} \frac{\partial}{\partial s} \frac{\partial}{\partial t} \mathbf{N} &= -\frac{\partial}{\partial s} [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \mathbf{T} + \frac{\partial}{\partial s} [\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] \\ &\quad + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]] \mathbf{B} - [\kappa [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \\ &\quad + \tau [\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]]] \mathbf{N}. \end{aligned}$$

On the other hand, we have

$$\begin{aligned} \frac{\partial}{\partial t} \frac{\partial}{\partial s} \mathbf{N} &= -\frac{\partial}{\partial t} \kappa \mathbf{T} + \frac{\partial}{\partial t} \tau \mathbf{B} \\ &\quad - \kappa [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau] \mathbf{N} - \kappa [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] \mathbf{B} + \tau \frac{\partial}{\partial t} \mathbf{B} \end{aligned}$$

Also,

$$\frac{\partial \tau}{\partial t} = - \langle \frac{\partial}{\partial t} \frac{\partial}{\partial s} \mathbf{B}, \mathbf{N} \rangle - \langle \frac{\partial}{\partial s} \mathbf{B}, \frac{\partial}{\partial t} \mathbf{N} \rangle.$$

From definition of flow, we have

$$\langle \mathbf{B}, \frac{\partial}{\partial t} \mathbf{B} \rangle = 0.$$

Thus, we obtain the theorem. The proof of theorem is completed. \square

Theorem 3.7. Let $\frac{\partial \alpha}{\partial t}$ be inextensible. Then,

$$\begin{aligned} \frac{\partial \mathbf{B}}{\partial t} &= \frac{1}{\tau} [\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]] \mathbf{T} \\ &\quad - [[\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]]] \mathbf{N}, \end{aligned} \quad (3.9)$$

where $\mathfrak{f}_1, \mathfrak{f}_2, \mathfrak{f}_3$ are smooth functions of time and arc length.

Proof: Using Frenet equations, we have

$$\begin{aligned} \frac{\partial \mathbf{B}}{\partial t} &= \frac{1}{\tau} [\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]] \mathbf{T} + \frac{1}{\tau} [\kappa [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] \\ &\quad - \frac{\partial \tau}{\partial t} + \frac{\partial}{\partial s} [\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]]] \mathbf{B} \\ &\quad - \frac{1}{\tau} [\tau [\frac{1}{\kappa} [\frac{\partial}{\partial s} [\frac{\partial \mathfrak{f}_3}{\partial s} + \mathfrak{f}_2 \tau] + \tau [\mathfrak{f}_1 \kappa + \frac{\partial \mathfrak{f}_2}{\partial s} - \mathfrak{f}_3 \tau]]]] \mathbf{N}. \end{aligned} \quad (3.10)$$

From definition of flow, we have

$$\langle \mathbf{B}, \frac{\partial}{\partial t} \mathbf{B} \rangle = 0. \quad (3.11)$$

Substituting (3.11) to (3.10), we have (3.9) as desired. This completes the proof. \square

Theorem 3.8. Let $\frac{\partial \alpha}{\partial t}$ be inextensible. Then,

$$\begin{aligned} & \kappa \left[\frac{1}{\tau} \left[\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \\ &= 2 \frac{\partial}{\partial s} \left[\left[\frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right], \end{aligned}$$

where f_1, f_2, f_3 are smooth functions of time and arc length.

Proof: Differentiating Eq. (3.9) with respect to s ,

$$\begin{aligned} \frac{\partial}{\partial s} \frac{\partial \mathbf{B}}{\partial t} &= \left[\frac{\partial}{\partial s} \left[\frac{1}{\tau} \left[\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right. \\ &\quad + \kappa \left[\left[\frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right] \mathbf{T} \\ &\quad + \left[\kappa \left[\frac{1}{\tau} \left[\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right. \\ &\quad \left. - \frac{\partial}{\partial s} \left[\left[\frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right] \right] \mathbf{N} \\ &\quad \left. - \tau \left[\left[\frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right] \right] \mathbf{B}. \end{aligned}$$

Also,

$$\begin{aligned} \frac{\partial}{\partial t} \frac{\partial \mathbf{B}}{\partial s} &= -\frac{\partial}{\partial t} \tau \mathbf{N} - \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \mathbf{T} \\ &\quad - \tau \frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \mathbf{B}. \end{aligned}$$

Hence, the proof is complete. \square

Corollary 3.9.

$$\begin{aligned} -\tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] &= \frac{\partial}{\partial s} \left[\frac{1}{\tau} \left[\frac{\partial \kappa}{\partial t} - \frac{\partial}{\partial s} \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \\ &\quad + \kappa \left[\left[\frac{1}{\kappa} \left[\frac{\partial}{\partial s} \left[\frac{\partial f_3}{\partial s} + f_2 \tau \right] + \tau \left[f_1 \kappa + \frac{\partial f_2}{\partial s} - f_3 \tau \right] \right] \right] \right], \end{aligned}$$

where f_1, f_2, f_3 are smooth functions of time and arc length.

Proof: It is obvious from Theorem (3.8). \square

Example 3.10. The helix is parametrized by

$$\gamma(u, t) = (A(t) \cos(u), A(t) \sin(u), B(t)u), \quad (3.12)$$

where A, B are functions only of time. The arc-length derivative is

$$(A^2 + B^2)^{\frac{1}{2}} \frac{\partial}{\partial s} = \frac{\partial}{\partial u}$$

and the evolution of γ is explicitly given by

$$\left(\frac{\partial A}{\partial t} \cos(u), \frac{\partial A}{\partial t} \sin(u), \frac{\partial B}{\partial t} u\right) = \frac{1}{A^2 + B^2} (-A \cos(u), -A \sin(u), 0).$$

Hence

$$\frac{\partial A}{\partial t} = -A (A^2 + B^2)^{-1}, \quad \frac{\partial B}{\partial t} = 0$$

and solutions are given by

$$\frac{A(t)^2}{2} + B^2 \log(A(t)) = -t + \frac{A(0)^2}{2} + B^2 \log(A(0)).$$

Note that, for positive B , $A(t)$ converges to, but never reaches, zero.

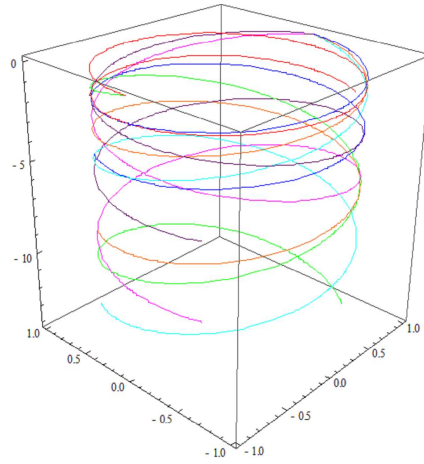


Figure 1: The equation Eq. (3.12) is illustrated colour Red, Blue, Purple, Orange, Magenta, Cyan, Green at the time $t = 1, t = 1.2, t = 1.4, t = 1.6, t = 1.8, t = 2, t = 2.2$, respectively.

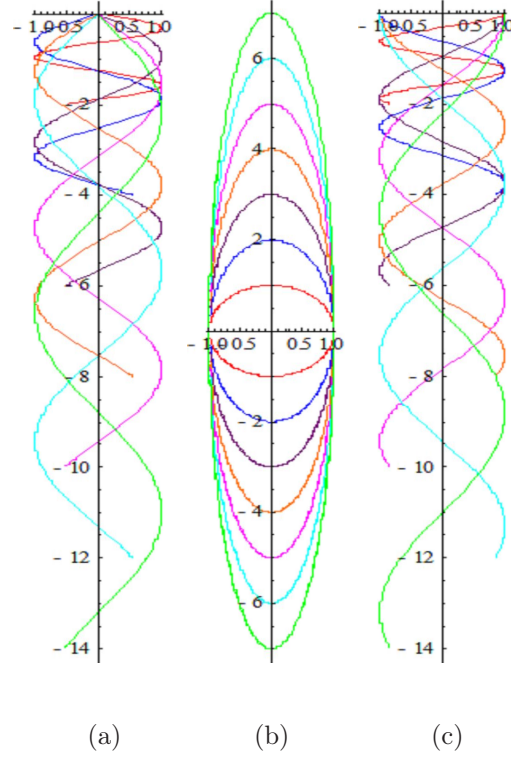


Figure 2: Projections of γ to yz , xy , xz planes are illustrated in (a), (b), (c), respectively.

References

1. U. Abresch, J. Langer: *The normalized curve shortening flow and homothetic solutions*, J. Differential Geom. 23 (1986), 175-196.
2. S.J. Altschuler and M.A. Grayson: *Shortening space curves and flow through singularities*, IMA preprint series (823), 1991.
3. B. Andrews: *Evolving convex curves*, Calculus of Variations and Partial Differential Equations, 7 (1998), 315-371.
4. J. H. Caltenco, R. Linares, and J. L. López-Bonilla: *Intrinsic geometry of curves and the Lorentz equation*, Czech. J. Phys. 52 (2002), 839-842.
5. M. Carmeli: *Motion of a charge in a gravitational field*, Phys. Rev. B 138 (1965), 1003-1007.
6. U. Ciftci: *A generalization of Lancert's theorem*, J. Geom. Phys. 59 (2009) 1597-1603.
7. P. Crouch, F. Silva Leite: *The dynamic interpolation problem: on Riemannian manifolds, Lie groups, and symmetric spaces*, J. Dynam. Control Systems 1 (2) (1995) 177-202.
8. M. P. Do Carmo: *Differential Geometry of Curves and Surfaces*, Prentice Hall, Englewood Cliffs, NJ, 1976.
9. A. Einstein: *Relativity: The Special and General Theory*, (New York: Henry Holt, 1920)
10. M. Gage, R.S. Hamilton: *The heat equation shrinking convex plane curves*, J. Differential Geom. 23 (1986), 69-96.

11. T. Körpınar, E. Turhan, V. Asil: *Tangent Bishop spherical images of a biharmonic B-slant helix in the Heisenberg group $Heis^3$* , Iranian Journal of Science & Technology, A (4) (2011), 265-271.
12. T. Körpınar and E. Turhan: *Biharmonic S-Curves According to Sabban Frame in Heisenberg Group $Heis^3$* , Bol. Soc. Paran. Mat., 31 (1) (2013), 205–211.
13. T. Körpınar, E. Turhan: *Parallel Surfaces to Normal Ruled Surfaces of General Helices in the Sol Space Sol^3* , Bol. Soc. Paran. Mat., (3s.) v. 31 2 (2013): 245–253.
14. D. Y. Kwon, F.C. Park: *Evolution of inelastic plane curves*, Appl. Math. Lett. 12 (1999), 115-119.
15. DY. Kwon , FC. Park, DP Chi: *Inextensible flows of curves and developable surfaces*, Appl. Math. Lett. 18 (2005), 1156-1162.
16. H.Q. Lu, J.S. Tdhunter, T.W. Sze: *Congruence conditions for nonplanar developable surfaces and their application to surface recognition*, CVGIP, Image Underst. 56 (1993), 265–285.
17. L. Ma and D. Chen : *Curve shortening in Riemannian manifold*, (preprint)
18. J. Milnor: *Curvatures of Left-Invariant Metrics on Lie Groups*, Advances in Mathematics 21 (1976), 293-329.
19. Z.-S. Shea, X. Chena and F. Hussaina: *Lie-group derivation of a multi-layer mixing length formula for turbulent channel and pipe flows*, arXiv:1112.6312
20. E. Turhan and T. Körpınar: *Parametric equations of general helices in the sol space Sol^3* , Bol. Soc. Paran. Mat. 31 (1) (2013), 99–104.
21. E. Turhan, T. Körpınar: *On Characterization Of Timelike Horizontal Biharmonic Curves In The Lorentzian Heisenberg Group $Heis^3$* , Zeitschrift für Naturforschung A- A Journal of Physical Sciences, 65a (2010), 641-648.
22. E. Turhan, T. Körpınar: *On Characterization Canal Surfaces around Timelike Horizontal Biharmonic Curves in Lorentzian Heisenberg Group $Heis^3$* , Z. Naturforsch. 66a (2011), 441 - 449
23. D.C. Wilcox: *Turbulence Modeling for CFD*. DCW Industries, 2006.

Talat Körpınar
 Muş Alparslan University, Department of Mathematics
 49250, Muş, Turkey
 E-mail address: talatkorpınar@gmail.com

and

Essin Turhan
 Fırat University, Department of Mathematics
 23119, Elazığ, Turkey
 E-mail address: essin.turhan@gmail.com