



Existence result of entropy solution for nonlinear elliptic problem without monotonicity condition in generalized sobolev spaces

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ABSTRACT: In the present paper we prove some existence results of entropy solution for nonlinear degenerate elliptic problems of the form $B(v) + H(x, v_n) = f$, in Musielak-Orlicz-Sobolev spaces, where $B(v) = -\operatorname{div}(b(x, v, \nabla v))$ is a Leray-Lions, operator defined from the musielak-Orlicz-sobolev spaces $W_0^1 L_\Psi(\Omega)$ into its dual $f \in L^1(\Omega)$, and no monotonicity strict condition is assumed on the function $b(x, s, \xi)$. The tool we use to overcome this difficulty is to investigate some techniques introduced by Minty's lemma.

Key Words: Elliptic problem, Entropy solutions, Musielak-Orlicz-Sobolev spaces, Compact imbedding, Δ_2 -condition..

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1. Introduction

The purpose of this note, is to show an existence of entropy solutions for some nonlinear Dirichlet problem as

$$\begin{cases} B(v) + H(x, v_n) = f(x) & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega, \end{cases} \quad (1.1)$$

where $B(v) = -\operatorname{div}(b(x, v, \nabla v))$, Ω is a bounded domain of \mathbb{R}^N , $N \geq 2$. $b(x, v, \nabla v) = (b_i(x, v, \nabla v))_{1 \leq i \leq N}$, $b_i : \Omega \times \mathbb{R} \times \mathbb{R}^N \rightarrow \mathbb{R}$ is a Carathéodory functions such that for all ξ, ξ' in \mathbb{R}^N ,

$$|b_i(x, s, \xi)| \leq |\phi_i(x)| + K_i \bar{P}^{-1}(\Psi(x, c_2|s|)) + K_i (\bar{\Psi}^{-1}\Psi(x, c_1|\xi|)), \quad (1.2)$$

$$(b(x, s, \xi) - b(x, s, \xi'))(\xi - \xi') \geq 0, \quad (1.3)$$

$$b(x, s, \xi) \xi \geq \Psi(x, \lambda_1|\xi|), \quad (1.4)$$

with $c_1, c_2, \lambda_1, K_i > 0$. Ψ, P be two Musielak functions such that $P \prec\prec \Psi$. Furthermore $\bar{\Psi}, \bar{P}$ be the conjugate functions of Ψ and P respectively, and $\phi_i \in E_{\bar{\Psi}}(\Omega)$.

Furthermore, the nonlinear term $H(x, s)$ verify

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$$H(x, s)s \geq 0, \quad (1.5)$$

$$\sup_{|s| \leq n} |H(x, s)| = h_n(x) \in L^1(\Omega), \quad (1.6)$$

$$f \in L^1(\Omega), \quad (1.7)$$

The Nonlinear elliptic problem as

$$\begin{cases} -\operatorname{div} b(x, v, \nabla v) = f - \operatorname{div} F & \text{in } \Omega \\ v = 0 & \text{on } \partial\Omega \end{cases}$$

with $f \in L^1(\Omega)$ and $F \in \left(L^{p'}(\Omega)\right)^N$, have been treated by Orsina et al. ([11]), by using the approach Minty's Lemma. Furthermore Badr EL HAJI et al. in their paper ([15]) and in weighted Orlicz-Sobolev spaces, they establish that the problem

$$\begin{cases} -\operatorname{div} b(x, v, \nabla v) = f(x) & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega, \end{cases}$$

with $f \in L^1$. admits at least one solution by using the idea of Minty's lemma, on the others hand and by the similar proof Rhoudaf et al ([1]) have shown the entropy solutions for the problem looks like

$$\begin{cases} -\operatorname{div} b(x, v, \nabla v) = h & \text{in } \Omega \\ v = 0 & \text{on } \partial\Omega. \end{cases}$$

with $h \in W^{-1,p'}(\Omega, \omega^*)$,

The mathematical researches dealing the existence of solutions to some problems parabolic and elliptic under a different assumptions is massive; we refer the reader to [3,4,5,6,7,13,14,15,16,17].

Our aim in this note is to solve the problem (1.1)(existence results) where the function $a_i(x, s, \xi)$ satisfy the large monotonicity condition and without adopting the almost everywhere convergence of the gradients, and in order to overcome this difficulties, we exploit the technique of Minty's lemma for proving the existence of an entropy solutions, However the approach that we used in the proof differs from that adopted by A. Benkirane et al used in [8]

The outline of this note is as follows. After giving the definition and some auxiliary results on anisotropic Sobolev space, and after recalling some essential assumptions which are necessary to have an existence solution, finally section 3 will be devoted to give our main results and their proofs.

2. Background

Here we give just some important definitions about Musielak-Orlicz spaces, for more details on this new space see ([22]).

2.1. Musielak-Orlicz-Sobolev spaces

Let Ω be an open subset of \mathbb{R}^n .

Definition 2.1 A Musielak-Orlicz function Ψ is a real-valued function defined in $\Omega \times \mathbb{R}_+$ such that
a) $\Psi(x, t)$ is an N -function i.e. convex, nondecreasing, continuous, $\Psi(x, 0) = 0, \Psi(x, t) > 0$ for all $t > 0$ and

$$\limsup_{t \rightarrow 0} \sup_{x \in \Omega} \frac{\Psi(x, t)}{t} = 0, \quad \liminf_{t \rightarrow \infty} \inf_{x \in \Omega} \frac{\Psi(x, t)}{t} = 0.$$

b) $\Psi(\cdot, t)$ is a Lebesgue measurable function.

For a Musielak-Orlicz function Ψ and a measurable function $v : \Omega \rightarrow \mathbb{R}$, we put

$$\rho_{\Psi, \Omega}(v) = \int_{\Omega} \Psi(x, |v(x)|) dx.$$

The set $K_{\Psi}(\Omega) = \{v : \Omega \rightarrow \mathbb{R} \text{ measurable} / \rho_{\Psi, \Omega}(v) < \infty\}$ is named the Musielak-Orlicz class. The Musielak-Orlicz space $L_{\Psi}(\Omega)$ is the vector space generated by $K_{\Psi}(\Omega)$, that is, $L_{\Psi}(\Omega)$ is the smallest linear space containing the set $K_{\Psi}(\Omega)$. That's to say

$$L_{\Psi}(\Omega) = \left\{ v : \Omega \rightarrow \mathbb{R} \text{ measurable} / \rho_{\Psi, \Omega} \left(\frac{v}{\lambda} \right) < \infty, \text{ for some } \lambda > 0 \right\}.$$

For a Musielak-Orlicz function Ψ we put: $\Psi(x, s) = \sup_{t>0} \{st - \Psi(x, t)\}$, Ψ is the conjugate Musielak-Orlicz function of Ψ in the sens of Young with respect to the variable s in the space $L_{\Psi}(\Omega)$

Let us define the two following equivalent norms as follows:

$$\|v\|_{\Psi, \Omega} = \inf \left\{ \lambda > 0 / \int_{\Omega} \Psi \left(x, \frac{|v(x)|}{\lambda} \right) dx \leq 1 \right\}, \text{ (The Luxemburg norm)}$$

$$\|v\|_{\Psi, \Omega} = \sup_{\|v\|_{\Psi} \leq 1} \int_{\Omega} |v(x)v(x)| dx, \text{ So-called Orlicz norm}$$

where Ψ is the Musielak Orlicz function conjugate to Ψ .

We will need the space $E_{\Psi}(\Omega)$ given by

$$E_{\Psi}(\Omega) = \left\{ w : \Omega \rightarrow \mathbb{R} \text{ measurable} / \rho_{\Psi, \Omega} \left(\frac{w}{\lambda} \right) < \infty, \text{ for all } \lambda > 0 \right\}.$$

A Musielak function Ψ is locally integrable on Ω if $\rho_{\Psi}(t\chi_D) < \infty$ for all $t > 0$ and all measurable $D \subset \Omega$ with $\text{meas}(D) < \infty$. Let Ψ a Musielak function locally integrable. Consequently $E_{\Psi}(\Omega)$ is separable ([22], Theorem 7.10).

We say that sequence of functions $v_n \in L_{\Psi}(\Omega)$ is modular convergent to $v \in L_{\Psi}(\Omega)$ if there exists $\lambda > 0$ such that

$$\lim_{n \rightarrow \infty} \rho_{\Psi, \Omega} \left(\frac{v_n - v}{\lambda} \right) = 0.$$

2.2. Auxilary lemmas

Lemma 2.1 [2]

If

$$\bar{\Psi}(x, 1) \leq c_1 \quad \text{a.e in } \Omega \text{ for some } c_1 > 0 \quad (2.1)$$

We said that Ψ verify the log-Hölder continuity where Ω be a bounded Lipschitz domain in \mathbb{R}^N ($N \geq 2$). Furthermore, $\mathfrak{D}(\Omega)$ is dense in $L_{\Psi}(\Omega)$ and in $W_0^1 L_{\Psi}(\Omega)$ for the modular convergence.

Lemma 2.2 ([23]) Let $F : \mathbb{R} \rightarrow \mathbb{R}$ be uniformly Lipschitzian, with $F(0) = 0$. Let Ψ be a Musielak-Orlicz function and let $\nu \in W_0^1 L_{\Psi}(\Omega)$. Then $F(\nu) \in W_0^1 L_{\Psi}(\Omega)$.

Moreover, if the set D of discontinuity points of F' is finite, we have

$$\frac{\partial}{\partial x_i} F(\nu) = \begin{cases} F'(\nu) \frac{\partial \nu}{\partial x_i} & \text{a.e in } \{x \in \Omega : \nu(x) \in D\} \\ 0 & \text{a.e in } \{x \in \Omega : \nu(x) \notin D\} \end{cases}$$

Lemma 2.3 [9] Assume that Ω verify the segment property and consider $\nu \in W_0^1 L_{\Psi}(\Omega)$. Consequently, there exists $(\nu_n) \subset \mathfrak{D}(\Omega)$:

$$\nu_n \rightarrow \nu \text{ in } W_0^1 L_{\Psi}(\Omega).$$

Moreover, if $\nu \in W_0^1 L_{\Psi}(\Omega) \cap L^{\infty}(\Omega)$ then $\|\nu_n\|_{\infty} \leq (N+1)\|\nu\|_{\infty}$.

Lemma 2.4 [21] Let (h_n) , $h \in L^1(\Omega)$ satisfying

i) $h_n \geq 0$ a.e in Ω ,

ii) $h_n \rightarrow h$ a.e in Ω ,

iii) $\int_{\Omega} h_n(x) dx \rightarrow \int_{\Omega} h(x) dx$.

Then $h_n \rightarrow h$ strongly in $L^1(\Omega)$.

Lemma 2.5 [10] Assume that $h_n \in L_{\Psi}(\Omega)$ converges in measure to h and if h_n bounded in $L_{\Psi}(\Omega)$, then $h \in L_{\Psi}(\Omega)$ and $h_n \rightarrow h$ for $\sigma(\Pi L_{\Psi}, \Pi E_{\Psi})$.

Lemma 2.6 [10] Suppose that $\nu_n, \nu \in L_{\Psi}(\Omega)$. If $\nu_n \rightarrow \nu$ for the modular convergence, then $\nu_n \rightarrow \nu$ for $\sigma(L_{\Psi}(\Omega), L_{\Psi}(\Omega))$.

Lemma 2.7 [18] If $P \prec \Psi$ and $\nu_n \rightarrow \nu$ for the modular convergence in $L_{\Psi}(\Omega)$ then $\nu_n \rightarrow \nu \in P(\Omega)$.

Now let us consider the truncation function T_k defined as follows

$$k \geq 0, T_k(s) = \max(-k, \min(k, s)).$$

3. Main Existence Theorem

Let Y be a closed subspace of $W^1 L_{\Psi}(\Omega)$ for $\sigma(\Pi L_{\Psi}, \Pi E_{\overline{\Psi}})$ and let

$$Y_0 = Y \cap W^1 L_{\Psi}(\Omega),$$

such that Y is the closure of Y_0 for $\sigma(\Pi L_{\Psi}, \Pi E_{\overline{\Psi}})$. In the next, we consider the complementary system (Y, Y_0, Z, Z_0) generated by Y i.e. Y_0^* can be identified to Z and Z_0^* can be identified to Y by the means $\langle \cdot, \cdot \rangle$. Let the mapping T (associated to the operator A) defined from

$$D(T) = \{v \in Y, b_0(x, v, \nabla v) \in L_{\overline{\Psi}}(\Omega), b_i(x, v, \nabla v) \in L_{\overline{\Psi}}(\Omega)\}$$

into Z by the formula

$$b(v, v) = \int_{\Omega} b_0(x, v, \nabla v) v(x) dx + \sum_{1 \leq i \leq N} \int_{\Omega} b_i(x, v, \nabla v) \frac{\partial v(x)}{\partial x_i} dx \quad \forall v \in Y_0.$$

We consider the complementary system

$$(Y, Y_0, Z, Z_0) = (W_0^1 L_{\Psi}(\Omega), W_0^1 E_{\Psi}(\Omega), W^{-1} E_{\overline{\Psi}}(\Omega), W^{-1} L_{\overline{\Psi}}(\Omega)).$$

Let us give the notion of entropy solution as follows

Definition 3.1 A function v (measurable) is named an entropy solution of the problem (1.1) if $T_k(v) \in W_0^1 L_{\Psi}(\Omega)$ for every $k > 0$ and such that

$$\int_{\Omega} b(x, v, \nabla v) \nabla T_k(v - \Theta) dx + \int_{\Omega} H(x, v_n) T_k(v - \Theta) dx \leq \int_{\Omega} f T_k(v - \Theta) dx$$

for every $\Theta \in W_0^1 E_{\Psi}(\Omega) \cap L^{\infty}(\Omega)$.

The existence results are stated as follows .

Theorem 3.1 If the hypotheses (1.2)-(1.7) holds, the problem (1.1) admit at least an entropy solution in the sense of the definition 3.1 .

3.1. Main Lemma

Lemma 3.1 *Suppose that v be a measurable function with $T_k(v) \in W_0^1 L_\Psi(\Omega)$ for every $k > 0$. Then*

$$\int_{\Omega} b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx + \int_{\Omega} H(x, v_n) T_k(v - \Theta) dx \leq \int_{\Omega} f T_k(v - \Theta) dx, \quad (3.1)$$

is equivalent to

$$\int_{\Omega} b(x, v, \nabla v) \nabla T_k(v - \Theta) dx + \int_{\Omega} H(x, v_n) T_k(v - \Theta) dx = \int_{\Omega} f T_k(v - \Theta) dx, \quad (3.2)$$

for every $\Theta \in W_0^1 L_\Psi(\Omega)$, and for every $k > 0$.

Proof In fact (3.2) implies (3.1) is easily proved adding and subtracting

$$\int_{\Omega} b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx,$$

and by thanking to assumption (1.3). Thus, it suffices to prove that (3.1) implies (3.2). Let $h, k > 0$, let us consider that $\lambda \in]-1, 1[$ and $\vartheta \in W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$. By taking, $\Theta = T_h(v - \lambda T_k(v - \vartheta)) \in W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$ as test in (3.1), we may obtain:

$$I_{hk} \leq J_{hk}, \quad (3.3)$$

where

$$\begin{aligned} I_{hk} &= \int_{\Omega} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta))) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx \\ &\quad + \int_{\Omega} H(x, v_n) T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx = I'_{hk} + I''_{hkk} \end{aligned}$$

and

$$J_{hk} = \int_{\Omega} f T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx.$$

Set

$$A_{hk} = \{x \in \Omega, |v - T_h(v - \lambda T_k(v - \vartheta))| \leq k\},$$

and

$$B_{hk} = \{x \in \Omega, |v - \lambda T_k(v - \vartheta)| \leq h\}.$$

Then, we get

$$\begin{aligned} I'_{hk} &= \int_{A_{kh} \cap B_{hk}} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta))) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx \\ &\quad + \int_{A_{kh} \cap B_{hk}^C} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta))) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx \\ &\quad + \int_{A_{kh}^C} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta))) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx. \end{aligned}$$

Since $\nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) \neq 0$ on A_{kh} , one has

$$\int_{A_{kh}^C} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta))) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta))) dx = 0. \quad (3.4)$$

Furthermore, if $x \in B_{hk}^C$, we have $\nabla T_h(v - \lambda T_k(v - \vartheta),) = 0$ and by according to (1.4), we conclude that,

$$\begin{aligned} & \int_{A_{kh} \cap B_{hk}^C} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx \\ &= \int_{A_{kh} \cap B_{hk}^C} b(x, v_n, 0) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx = 0. \end{aligned} \quad (3.5)$$

According to (3.4) and (3.5), one has

$$I'_{hk} = \int_{A_{kh} \cap B_{hk}} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx.$$

Letting $h \rightarrow +\infty, |\lambda| \leq 1$, we obtain

$$A_{kh} \rightarrow \{x, |\lambda| |T_k(v - \vartheta)| \leq h\} = \Omega, \quad (3.6)$$

$$B_{hk} \rightarrow \Omega \quad \text{which gives} \quad A_{kh} \cap B_{hk} \rightarrow \Omega. \quad (3.7)$$

By applying Lebesgue theorem, we deduce that

$$\begin{aligned} & \lim_{h \rightarrow +\infty} \int_{A_{kh} \cap B_{hk}} b(x, v_n, \nabla T_h(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx \\ &= \lambda \int_{\Omega} b(x, v_n, \nabla(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - \vartheta) \, dx. \end{aligned} \quad (3.8)$$

Thus,

$$\lim_{h \rightarrow +\infty} I'_{hk} = \lambda \int_{\Omega} b(x, v_n, \nabla(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - \vartheta) \, dx.$$

Furthermore we can easily remark that

$$\lim_{h \rightarrow +\infty} \int_{\Omega} H(x, v_n) T_k(v - T_h(v - \lambda T_k(v - \vartheta))) \, dx = \lambda \int_{\Omega} H(x, v_n) T_k[v - \vartheta] \, dx$$

consequently

$$\lim_{h \rightarrow +\infty} I_{hk} = \lambda \int_{\Omega} b(x, v_n, \nabla(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - \vartheta) \, dx + \lambda \int_{\Omega} H(x, v_n) T_k[v - \vartheta] \, dx. \quad (3.9)$$

On the other hand, we have,

$$J_{hk} = \int_{\Omega} f T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx.$$

Then

$$\lim_{h \rightarrow +\infty} \int_{\Omega} f T_k(v - T_h(v - \lambda T_k(v - \vartheta),)) \, dx = \lambda \int_{\Omega} f T_k(v - \vartheta) \, dx, \quad (3.10)$$

i.e.,

$$\lim_{h \rightarrow +\infty} J_{hk} = \lambda \int_{\Omega} f T_k(v - \vartheta) \, dx. \quad (3.11)$$

Tanking to (3.9), (3.11) and passing to the limit in (3.3), one has,

$$\lambda \left(\int_{\Omega} b(x, v_n, \nabla(v - \lambda T_k(v - \vartheta),)) \nabla T_k(v - \vartheta) \, dx + \int_{\Omega} H(x, v_n) T_k[v - \vartheta] \, dx \right) \leq \lambda \left(\int_{\Omega} f T_k(v - \vartheta) \, dx \right)$$

for every $\vartheta \in W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$, and for every $k > 0$. Taking $\lambda > 0$ dividing by λ , and then letting $\lambda \rightarrow 0$, we get

$$\int_{\Omega} b(x, v, \nabla v) \nabla T_k(v - \vartheta) dx + \int_{\Omega} H(x, v_n) T_k[v - \vartheta] dx \leq \int_{\Omega} f T_k(v - \vartheta) dx. \quad (3.12)$$

for $\lambda < 0$, dividing by λ , and then letting $\lambda \rightarrow 0$, we have

$$\int_{\Omega} b(x, v, \nabla v) \nabla T_k(v - \vartheta) dx + \int_{\Omega} H(x, v_n) T_k[v - \vartheta] dx \geq \int_{\Omega} f T_k(v - \vartheta) dx. \quad (3.13)$$

According (3.12) and (3.13), we deduce that :

$$\int_{\Omega} b(x, v, \nabla v) \nabla T_k(v - \vartheta) dx + \int_{\Omega} H(x, v_n) T_k[v - \vartheta] dx = \int_{\Omega} f T_k(v - \vartheta) dx. \quad (3.14)$$

This completes the proof of Lemma 3.1.

3.2. Proof of Theorem 3.1

3.2.1. Approximate problem and a priori estimate. For $n \in \mathbb{N}$, define $f_n := T_n(f)$. Let v_n be solution in $W_0^1 L_\Psi(\Omega)$ of the problem

$$\begin{cases} -\operatorname{div}(b(x, v_n, \nabla v_n)) + H_n(x, v_n) = f_n & \text{in } \Omega \\ v_n = 0 & \text{on } \partial\Omega, \end{cases} \quad (3.15)$$

where,

$$H_n(x, s) = \frac{H(x, s)}{1 + \frac{1}{n}|H(x, s)|},$$

which exists thanks to [20]. Choosing $T_k(v_n)$ as test function in (3.15), we have

$$\int_{\Omega} b(x, v_n, \nabla v_n) \nabla T_k(v_n) dx + \int_{\Omega} H_n(x, v_n) T_k(v_n) dx = \int_{\Omega} f_n T_k(v_n) dx,$$

Using $\nabla T_k(v_n) = \nabla v_n \chi_{\{|v_n| \leq k\}}$ and in view of (1.4), then since $H_n(x, v_n) T_k(v_n) \geq 0$, we obtain

$$\int_{\Omega} b(x, v_n, \nabla v_n) \nabla T_k(v_n) dx \geq \int_{\Omega} \Psi(x, \lambda_1 |\nabla T_k(v_n)|) dx,$$

then

$$\int_{\Omega} \Psi(x, \lambda_1 |\nabla T_k(v_n)|) dx \leq k \|f\|_{L^1(\Omega)}. \quad (3.16)$$

Then

$$\int_{\Omega} \Psi(x, \lambda_1 |\nabla T_k(v_n)|) dx \leq C_1 k, \quad (3.17)$$

where C_1 is a constant independently of n .

3.2.2. Locally convergence of v_n in measure. Taking $\frac{1}{\lambda} |T_k(v_n)|$ in (3.15) and using (3.17), one has

$$\int_{\Omega} \Psi(x, \lambda_1 \frac{|\nabla T_k(v_n)|}{\lambda}) dx \leq \int_{\Omega} \Psi(x, \lambda_1 |\nabla T_k(v_n)|) dx \leq C_1 k. \quad (3.18)$$

Then we deduce by using (3.18), that

$$\begin{aligned} \operatorname{meas}\{|v_n| > k\} &\leq \frac{1}{\inf_k \Psi(x, \frac{k}{\lambda})} \int_{\{|v_n| > k\}} \Psi(x, \frac{|v_n(x)|}{\lambda}) dx \\ &\leq \frac{1}{\inf_k \Psi(x, \frac{k}{\lambda})} \int_{\Omega} \Psi(x, \frac{1}{\lambda} |T_k(v_n)|) dx \\ &\leq \frac{C_1 k}{\inf_k \Psi(x, \frac{k}{\lambda})} \quad \forall n, \quad \forall k \geq 0. \end{aligned} \quad (3.19)$$

For any $\beta > 0$, we have

$$\text{meas}\{|v_n - v_m| > \beta\} \leq \text{meas}\{|v_n| > k\} + \text{meas}\{|v_m| > k\} + \text{meas}\{|T_k(v_n) - T_k(v_m)| > \beta\},$$

and so that

$$\text{meas}\{|v_n - v_m| > \beta\} \leq \frac{2C_1 k}{\inf_k \Psi(x, \frac{k}{\lambda})} + \text{meas}\{|T_k(v_n) - T_k(v_m)| > \beta\}. \quad (3.20)$$

According to (3.17) and by using Poincaré inequality in generalized -Sobolev spaces, we may obtain that $(T_k(v_n))$ is bounded in $W_0^1 L_\Psi(\Omega)$, and then there exists $\omega_k \in W_0^1 L_\Psi(\Omega)$ such that $T_k(v_n) \rightharpoonup \omega_k$ weakly in $W_0^1 L_\Psi(\Omega)$ for $\sigma(\Pi L_\Psi, \Pi E_{\bar{\Psi}})$; strongly in $E_{\bar{\Psi}}(\Omega)$ and a.e. in Ω .

Thus, we can consider that $(T_k(v_n))_n$ is a Cauchy sequence in measure in Ω .

Let $\varepsilon > 0$, then by (3.20) and in view of $\frac{2C_1 k}{\inf_k \Psi(x, \frac{k}{\lambda})} \rightarrow 0$ as $k \rightarrow +\infty$ there exists some $k = k(\varepsilon) > 0$ such that

$$\text{meas}\{|v_n - v_m| > \lambda\} < \varepsilon, \quad \text{for all } n, m \geq h_0(k(\varepsilon), \lambda).$$

Consequently, v_n is a Cauchy sequence in measure, thus, v_n converges almost everywhere to some measurable function v . Finally, there exist a subsequence of $\{v_n\}_n$, still indexed by n , and a function $v \in W_0^1 L_\Psi(\Omega)$ such that

$$\begin{cases} v_n u_n \rightharpoonup v & \text{weakly in } W_0^1 L_\Psi(\Omega) \text{ for } \sigma(\Pi L_\Psi, \Pi E_{\bar{\Psi}}) \\ v_n \rightarrow v & \text{strongly in } E_\Psi(\Omega) \text{ and a.e. in } \Omega. \end{cases} \quad (3.21)$$

3.2.3. Equi-integrability of nonlinearities. Here we try to prove that

$$H_n(x, v_n) \rightarrow H(x, v_n) \text{ strongly in } L^1(\Omega) \quad (3.22)$$

For this we choose $T_{l+1}(v_n) - T_l(v_n)$ as test in (3.15), we may get

$$\begin{aligned} \int_{\Omega} b(x, v_n, \nabla v_n) \nabla (T_{l+1}(v_n) - T_l(v_n)) dx + \int_{\Omega} H_n(x, v_n) (T_{l+1}(v_n) - T_l(v_n)) dx \\ = \int_{\Omega} f(T_{l+1}(v_n) - T_l(v_n)) dx \end{aligned}$$

which implies that

$$\begin{aligned} \int_{\{l \leq |v_n| \leq l+1\}} b(x, v_n, \nabla v_n), \nabla v_n dx + \int_{\{|v_n| \geq l+1\}} |H_n(x, v_n)| dx \\ \leq c \int_{\{|v_n| \geq l\}} |f| dx. \end{aligned}$$

Thus by (1.4), we have

$$\int_{\{|v_n| \geq l+1\}} |H_n(x, v_n)| dx \leq c \int_{\{|v_n| \geq l\}} |f_n| dx.$$

Let $\varepsilon > 0$, then there exist $l(\varepsilon) \geq 1$ such that

$$\int_{\{|v_n| > l(\varepsilon)\}} |H_n(x, v_n)| dx \leq \frac{\varepsilon}{2}, \quad (3.23)$$

For any measurable subset $E \subset \Omega$, we have

$$\begin{aligned} \int_E |H_n(x, v_n)| dx &\leq \int_{E \cap \{|v_n| \leq l(\varepsilon)\}} |H_n(x, v_n)| dx + \int_{E \cap \{|v_n| > l(\varepsilon)\}} |H_n(x, v_n)| dx \\ &\leq \int_E |h_{l(\varepsilon)}(x)| dx + \int_{E \cap \{|v_n| > l(\varepsilon)\}} |g_n(x, v_n)| dx. \end{aligned}$$

According to (1.6) there exist $\eta(\varepsilon) > 0$ such that

$$\int_E |h_{l(\varepsilon)}(x)| dx \leq \frac{\varepsilon}{2}, \quad (3.24)$$

for all E such that $\text{meas}(E) < \eta(\varepsilon)$

Finally, by combining (3.23) and (3.24) one easily has $\int_E |H_n(x, v_n)| dx \leq \varepsilon$, for all E such that $\text{meas}(E) < \eta(\varepsilon)$

3.2.4. An intermediate Inequality. In this step, we shall prove that for $\Theta \in W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$, we have

$$\int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx + \int_\Omega H_n(x, v_n) T_k(v_n - \Theta) dx \leq \int_\Omega f_n T_k(v_n - \Theta) dx. \quad (3.25)$$

Now, it suffice to take $T_k(v_n - \Theta)$ as test in (3.15), with Θ in $W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$, in order to get

$$\int_\Omega b(x, v_n, \nabla v_n) \nabla T_k(v_n - \Theta) dx + \int_\Omega H_n(x, v_n) T_k(v_n - \Theta) dx = \int_\Omega f_n T_k(v_n - \Theta) dx. \quad (3.26)$$

Adding and subtracting the term $\int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx$ i.e.,

$$\begin{aligned} & \int_\Omega b(x, v_n, \nabla v_n) \nabla T_k(v_n - \Theta) dx + \int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx \\ & - \int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx + \int_\Omega H_n(x, v_n) T_k(v_n - \Theta) dx = \int_\Omega f_n T_k(v_n - \Theta) dx. \end{aligned} \quad (3.27)$$

According to the hypothesis (1.3), we may obtain

$$\int_\Omega (b(x, v_n, \nabla v_n) - b(x, v_n, \nabla \Theta)) \nabla T_k(v_n - \Theta) dx \geq 0. \quad (3.28)$$

According (3.26) and (3.28), we obtain (3.25).

3.2.5. Passing to the limit. Here we show that,

$$\int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx + \int_\Omega H(x, v) T_k(v - \Theta) dx \leq \int_\Omega f T_k(v - \Theta) dx, \text{ for } \Theta \in W_0^1 L_\Psi(\Omega) \cap L^\infty(\Omega)$$

Firstly, we verify that

$$\int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx \rightarrow \int_\Omega b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx \text{ as } n \rightarrow +\infty.$$

Since $T_\mathbb{E}(v_n) \rightharpoonup T_\mathbb{E}(v)$ weakly in $W_0^1 L_\Psi(\Omega)$, with $\mathbb{E} = k + \|\Theta\|_\infty$, then

$$T_k(v_n - \Theta) \rightharpoonup T_k(v - \Theta) \text{ in } W_0^1 L_\Psi(\Omega), \quad (3.29)$$

this entails that

$$\frac{\partial T_k}{\partial x_i}(v_n - \Theta) \rightharpoonup \frac{\partial T_k}{\partial x_i}(v - \Theta) \text{ weakly in } L_\Psi(\Omega) \quad \forall i = 1, \dots, N. \quad (3.30)$$

Show that

$$b(x, T_\mathbb{E}(v_n), \nabla \Theta) \rightarrow b(x, T_\mathbb{E}(v), \nabla \Theta) \text{ strongly in } (L_{\overline{\Psi}}(\Omega))^N.$$

According to hypothesis (1.2), we may get

$$|b_i(x, T_\mathbb{E}(v_n), \nabla \Theta)| \leq |\Theta_i(x)| + K_i \overline{P}^{-1}(\Psi(x, c_2 |T_\mathbb{E}(v_n)|)) + K_i \overline{\Psi}^{-1}(\Psi(x, c_1 |\nabla \Theta|)),$$

with β and μ are positive constants. Since $T_{\mathbb{E}}(v_n) \rightharpoonup T_{\mathbb{E}}(v)$ weakly in $W_0^1 L_{\Psi}(\Omega)$ and $W_0^1 L_{\Psi}(\Omega) \hookrightarrow L_{\overline{\Psi}}(\Omega)$, then $T_{\mathbb{E}}(v_n) \rightarrow T_{\mathbb{E}}(v)$ strongly in $L_{\Psi}(\Omega)$ and a.e. in Ω , hence

$$|b(x, T_{\mathbb{E}}(v_n), \nabla \Theta)| \rightarrow |b(x, T_{\mathbb{E}}(v), \nabla \Theta)| \text{ a.e. in } \Omega.$$

and

$$|\Theta_i(x)| + K_i \overline{P}^{-1}(\Psi(x, c_2 |T_{\mathbb{E}}(v_n)|)) + K_i \overline{\Psi}^{-1}\Psi(x, c_1 |\nabla \Theta|) \rightarrow$$

$$|\Theta_i(x)| + K_i \overline{P}^{-1}(\Psi(x, c_2 |T_{\mathbb{E}}(v)|)) + K_i \overline{\Psi}^{-1}\Psi(x, c_1 |\nabla \Theta|),$$

a.e. in Ω . then, Vitali's theorem implies that

$$b(x, T_{\mathbb{E}}(v_n), \nabla \Theta) \rightarrow b(x, T_{\mathbb{E}}(v), \nabla \Theta) \text{ strongly in } (L_{\overline{\Psi}}(\Omega))^N, \text{ as } n \rightarrow \infty. \quad (3.31)$$

According (3.30) and (3.31), we obtain

$$\int_{\Omega} b(x, v_n, \nabla \Theta) \nabla T_k(v_n - \Theta) dx \rightarrow \int_{\Omega} b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx \text{ as } n \rightarrow +\infty. \quad (3.32)$$

Furthermore, we establish that

$$\int_{\Omega} f_n T_k(v_n - \Theta) dx \rightarrow \int_{\Omega} f T_k(v - \Theta) dx. \quad (3.33)$$

We have $f_n T_k(v_n - \Theta) \rightarrow f T_k(v - \Theta)$ a.e. in Ω and $|f T_k(v_n - \Theta)| \leq k|f|$,

By applying Vitali's theorem, we may obtain (3.33).

Similarly in view of (3.22) we may prove that

$$\int_{\Omega} H_n(x, v_n) T_k(v_n - \Theta) dx \rightarrow \int_{\Omega} H(x, v_n) T_k(v - \Theta) dx \text{ as } n \rightarrow \infty. \quad (3.34)$$

According to (3.32), (3.33) and (3.34) we may pass to the limit in (3.25), so that $\forall \Theta \in W_0^1 L_{\Psi}(\Omega) \cap L^{\infty}(\Omega)$, we deduce

$$\int_{\Omega} b(x, v_n, \nabla \Theta) \nabla T_k(v - \Theta) dx + \int_{\Omega} H(x, v_n) T_k(v - \Theta) dx \leq \int_{\Omega} f T_k(v - \Theta) dx.$$

According to the Lemma 3.1, we may deduce that v is an entropy solution of the problem (1.1). This achieve the proof of theorem 3.1.

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