



Geometry of η -Ricci Yamabe Soliton on Nearly Sasakian Manifold

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ABSTRACT: The present paper is devoted to study η -Ricci-Yamabe soliton on nearly-Sasakian manifolds. We examine Ricci-semisymmetry and Einstein-semisymmetric on nearly-sasakian manifold to obtain condition for shrinking or steady or expanding. Further, we analyze the 3-dimensional nearly-Sasakian manifolds admitting η -Ricci-Yamabe soliton satisfying certain geometric conditions.

Key Words: Nearly-Sasakian manifold, η -Ricci-Yamabe soliton, Einstein tensor, η -Einstein manifold, Projective curvature tensor.

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1. INTRODUCTION

A Ricci soliton is nothing but a generalised concept of the Einstein metric. In a Riemannian manifold $(\Omega, \mathbf{g})$ of dimension n , the metric $\mathbf{g}$ is called a Ricci soliton if

$$(\mathfrak{L}_{\Upsilon} \mathbf{g})(\mathcal{A}_1, \mathcal{A}_2) + 2S(\mathcal{A}_1, \mathcal{A}_2) + 2\lambda_1 \mathbf{g}(\mathcal{A}_1, \mathcal{A}_2) = 0, \quad (1.1)$$

where $\mathfrak{L}_{\Upsilon} \mathbf{g}$ is the Lie derivative of the Riemannian metric $\mathbf{g}$ along the vector field Υ , S is the Ricci tensor, λ_1 is a constant and $\mathcal{A}_1, \mathcal{A}_2$ are arbitrary vector fields on $T\Omega$. The Ricci soliton is said to be expanding if $\lambda_1 = \text{positive}$; shrinking if $\lambda_1 = \text{negative}$; and steady if $\lambda_1 = 0$. A Ricci soliton with $\Upsilon = 0$ reduces to an Einstein equation. Cho and Kimura(2009) studied an η -Ricci soliton $(\mathbf{g}, \Upsilon, \lambda_1, \lambda_2)$ as the generalization of the Ricci soliton $(\mathbf{g}, \Upsilon, \lambda_1)$ [4]. They have explained the extension of Ricci soliton as a notion of η -Ricci soliton [4]. Hamilton(1988) initiated the concepts – Ricci flow and Yamabe flow [5]. Some related works can be found in [1-3,6-17]. The η -Ricci Yamabe soliton (η -RYS) concept was defined by Siddiqi and Akyol [18] as follows:

$$(\mathfrak{L}_{\Upsilon} \mathbf{g})(\mathcal{A}, \mathcal{B}) + 2\rho S(\mathcal{A}, \mathcal{B}) + (2\lambda_1 - qr)\mathbf{g}(\mathcal{A}, \mathcal{B}) + 2\lambda_2 \eta(\mathcal{A})\eta(\mathcal{B}) = 0, \quad (1.2)$$

where λ_1 and λ_2 are constants.

In this paper, Ricci tensor for nearly Sasakian η -RYS is attained, and Ricci-semisymmetry and Einstein-semisymmetry conditions are studied. Nearly-Sasakian η -RYS satisfying the situation $S \cdot R =$

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0 is discussed. η -RYS on a 3-dimensional nearly-Sasakian manifold satisfying the situation $Q \cdot R = 0$ is discussed. Further, η -RYS on a 3-dimensional nearly-Sasakian manifold satisfying the situation $R(\mathcal{A}, \varsigma) \cdot P - P(\mathcal{A}, \varsigma) \cdot R = 0$ is discussed.

2. BASICS

Let Ω denotes a smooth manifold of dimension n . An almost contact structure on Ω is given by $(\Theta, \varsigma, \eta)$, where Θ is a one-one tensor field, ς is a vector field, and η is a 1-form on Ω satisfying the relations

$$\Theta^2 \mathcal{A} = -\mathcal{A} + \eta(\mathcal{A})\varsigma, \quad \eta(\varsigma) = 1, \quad \Theta(\varsigma) = 0, \quad \eta(\varsigma \mathcal{A}) = 0, \quad (2.1)$$

$$\mathfrak{g}(\Theta \mathcal{A}, \Theta \mathcal{B}) = \mathfrak{g}(\mathcal{A}, \mathcal{B}) - \eta(\mathcal{A})\eta(\mathcal{B}), \quad (2.2)$$

for all $\mathcal{A}, \mathcal{B}$ in the Lie algebra $\chi(\Omega)$ of all vector fields on Ω . An almost contact manifold provided with a Riemannian metric is called as an almost contact metric manifold(ACMM) and it is expressed by $\Omega(\Theta, \varsigma, \eta, \mathfrak{g})$. In addition, on a Nearly-Sasakian manifold(NSM), below mentioned identities are preserved:

$$R(\mathcal{A}_1, \mathcal{A}_2)\varsigma = \eta(\mathcal{A}_2)\mathcal{A}_1 - \eta(\mathcal{A}_1)\mathcal{A}_2, \quad (2.3)$$

$$R(\varsigma, \mathcal{A}_1)\mathcal{A}_2 = -R(\mathcal{A}_1, \varsigma)\mathcal{A}_2 = \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2)\varsigma - \eta(\mathcal{A}_2)\mathcal{A}_1, \quad (2.4)$$

$$S(\mathcal{A}_1, \varsigma) = (n-1)\eta(\mathcal{A}_1), \quad (2.5)$$

$$Q\varsigma = (n-1)\varsigma, \quad (2.6)$$

$$S(\Theta \mathcal{A}_1, \Theta \mathcal{A}_2) = S(\mathcal{A}_1, \mathcal{A}_2) - (n-1)\eta(\mathcal{A}_1)\eta(\mathcal{A}_2), \quad (2.7)$$

for all vector fields $\mathcal{A}_1, \mathcal{A}_2$ on Ω . The following definitions are needed:

Definition 2.1 Let Ω be an n -dimensional NSM. We say that Ω is Ricci semisymmetric if $R \cdot S = 0$.

Definition 2.2 Let Ω be an n -dimensional NSM. We say that Ω is Einstein semisymmetric if $R \cdot E = 0$, where E is the Einstein tensor given by

$$E(\mathcal{A}_1, \mathcal{A}_2) = S(\mathcal{A}_1, \mathcal{A}_2) - \left(\frac{r}{n}\right) \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2), \quad (2.8)$$

where S is the Ricci tensor and r is a scalar curvature of the manifold Ω .

Definition 2.3 We say that a manifold Ω is an η -Einstein manifold if its Ricci tensor is given by

$$S(\mathcal{A}_1, \mathcal{A}_2) = c_1 \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2) + c_2 \eta(\mathcal{A}_1)\eta(\mathcal{A}_2),$$

for all vector fields $\mathcal{A}_1, \mathcal{A}_2$ on Ω , where c_1 and c_2 are smooth functions on Ω .

3. η -RYS ON NSM

Let Ω be a NSM. Let $(\mathfrak{g}, \varsigma, \lambda_1, \lambda_2)$ be an η -RYS on Ω . Then by (1.2),

$$2\rho S(\mathcal{A}, \mathcal{B}) = -(\mathfrak{L}_v \mathfrak{g})(\mathcal{A}, \mathcal{B}) - (2\lambda_1 - qr)\mathfrak{g}(\mathcal{A}, \mathcal{B}) - 2\lambda_2 \eta(\mathcal{A})\eta(\mathcal{B}). \quad (3.1)$$

On a nearly-Sasakian manifold, from a nearly-Sasakian structure $\nabla_{\mathcal{A}} \varsigma = -\Theta \mathcal{A}$, we have

$$(\mathfrak{L}_v \mathfrak{g})(\mathcal{A}, \mathcal{B}) = 0. \quad (3.2)$$

Using (3.2) in (3.1), we obtain

$$S(\mathcal{A}, \mathcal{B}) = \frac{(qr - 2\lambda_1)}{2\rho} \mathfrak{g}(\mathcal{A}, \mathcal{B}) - \frac{\lambda_2}{\rho} \eta(\mathcal{A})\eta(\mathcal{B}), \quad (3.3)$$

$$Q\mathcal{A} = \frac{(qr - 2\lambda_1)}{2\rho} \mathcal{A} - \frac{\lambda_2}{\rho} \eta(\mathcal{A})\varsigma. \quad (3.4)$$

Putting $\mathcal{B} = \varsigma$ in (3.3) and (3.4), we obtain

$$S(\mathcal{A}, \varsigma) = \frac{(qr - 2\lambda_1 - 2\lambda_2)}{2\rho} \eta(\mathcal{A}), \quad (3.5)$$

$$Q_\varsigma = \frac{(qr - 2\lambda_1 - 2\lambda_2)}{2\rho} \varsigma. \quad (3.6)$$

Comparing (3.5) with (2.5), we obtain

$$n = 1 + \frac{(qr - 2\lambda_1 - 2\lambda_2)}{2\rho}.$$

Consequently, we obtain the following outcome.

Theorem 3.1 *For an η -RYS on a NSM Ω , the Ricci tensor is in the form (3.3) and $\lambda_1 + \lambda_2 = \rho(1-n) + \frac{qr}{2}$.*

4. η -RYS ON RICCI-SEMISSYMMETRIC NSM

Consider a Ricci-semisymmetric NSM Ω . By Definition 2.1,

$$S(R(\varsigma, \mathcal{A}_1)\mathcal{A}_2, \mathcal{A}_3) + S(\mathcal{A}_2, R(\varsigma, \mathcal{A}_1)\mathcal{A}_3) = 0. \quad (4.1)$$

Suppose Ricci-semisymmetric NSM admits an η -RYS. Then by Eq.(3.3), Eq.(4.1) becomes

$$\begin{aligned} & \left[\frac{qr - 2\lambda_1}{2\rho} \right] \mathfrak{g}(R(\varsigma, \mathcal{A}_1)\mathcal{A}_2, \mathcal{A}_3) - \frac{\lambda_2}{\rho} \eta(R(\varsigma, \mathcal{A}_1)\mathcal{A}_2) \eta(\mathcal{A}_3) \\ & + \left[\frac{qr - 2\lambda_1}{2\rho} \right] \mathfrak{g}(\mathcal{A}_2, R(\varsigma, \mathcal{A}_1)\mathcal{A}_3) - \frac{\lambda_2}{\rho} \eta(R(\varsigma, \mathcal{A}_1)\mathcal{A}_3) \eta(\mathcal{A}_2) = 0. \end{aligned} \quad (4.2)$$

Using (2.4), we have

$$\begin{aligned} & \left[\frac{qr - 2\lambda_1}{2\rho} \right] \mathfrak{g}(\mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2)\varsigma - \eta(\mathcal{A}_2)\mathcal{A}_1, \mathcal{A}_3) - \frac{\lambda_2}{\rho} \eta(\mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2)\varsigma - \eta(\mathcal{A}_2)\mathcal{A}_1) \eta(\mathcal{A}_3) + \\ & \left[\frac{qr - 2\lambda_1}{2\rho} \right] \mathfrak{g}(\mathcal{A}_2, \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_3)\varsigma - \eta(\mathcal{A}_3)\mathcal{A}_1) - \frac{\lambda_2}{\rho} \eta(\mathfrak{g}(\mathcal{A}_1, \mathcal{A}_3)\varsigma - \eta(\mathcal{A}_3)\mathcal{A}_1) \eta(\mathcal{A}_2) = 0 \end{aligned} \quad (4.3)$$

which implies

$$\frac{\lambda_2}{\rho} [\mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2) \eta(\mathcal{A}_3) + \eta(\mathcal{A}_2) \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_3) - 2\eta(\mathcal{A}_1) \eta(\mathcal{A}_2) \eta(\mathcal{A}_3)] = 0. \quad (4.4)$$

Putting $\mathcal{A}_3 = \varsigma$ in (4.4), yields

$$\frac{\lambda_2}{\rho} [\mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2) - \eta(\mathcal{A}_1) \eta(\mathcal{A}_2)] = 0. \quad (4.5)$$

Replace $\mathcal{A}_1$ by $\Theta\mathcal{A}_1$ and $\mathcal{A}_2$ by $\Theta\mathcal{A}_2$, we obtain

$$\frac{\lambda_2}{\rho} \mathfrak{g}(\Theta\mathcal{A}_1, \Theta\mathcal{A}_2) = 0. \quad (4.6)$$

Hence $\lambda_2 = 0$. Therefore, by Theorem 3.1, we have $\lambda_1 = \rho(1-n) + \frac{qr}{2}$. Thus, we have

Theorem 4.1 *A nearly-Sasakian η -RYS cannot be Ricci-semisymmetric and $\lambda_1 = \rho(1-n) + \frac{qr}{2}$.*

5. NEARLY-SASAKIAN η -RYS SATISFYING $S \cdot R = 0$

Consider a nearly-Sasakian η -RYS satisfying

$$(S(\mathcal{A}_1, \mathcal{A}_2) \cdot R)(\mathcal{U}, \Upsilon)\mathcal{A}_3 = 0, \quad (5.1)$$

for all $\mathcal{A}_1, \mathcal{A}_2, \mathcal{A}_3, \mathcal{U}$ and Υ on Ω . Equation (5.1) implies that

$$\begin{aligned} (\mathcal{A}_1 \Lambda_s \mathcal{A}_2)R(\mathcal{U}, \Upsilon)\mathcal{A}_3 + R((\mathcal{A}_1 \Lambda_s \mathcal{A}_2)\mathcal{U}, \Upsilon)\mathcal{A}_3 + R(\mathcal{U}, (\mathcal{A}_1 \Lambda_s \mathcal{A}_2)\Upsilon)\mathcal{A}_3 \\ + R(\mathcal{U}, \Upsilon)(\mathcal{A}_1 \Lambda_s \mathcal{A}_2)\mathcal{A}_3 = 0, \end{aligned} \quad (5.2)$$

where $(\mathcal{A}_1 \Lambda_s \mathcal{A}_2)\mathcal{A}_3 = S(\mathcal{A}_2, \mathcal{A}_3)\mathcal{A}_1 - S(\mathcal{A}_1, \mathcal{A}_3)\mathcal{A}_2$. Using (5.2) in (5.1) and taking $\mathcal{A}_1 = \varsigma$, we obtain

$$\begin{aligned} S(\mathcal{A}_2, R(\mathcal{U}, \Upsilon)\mathcal{A}_3)\varsigma - S(\varsigma, R(\mathcal{U}, \Upsilon)\mathcal{A}_3)\mathcal{A}_2 + S(\mathcal{A}_2, \mathcal{U})R(\varsigma, \Upsilon)\mathcal{A}_3 - S(\varsigma, \mathcal{U}), R(\mathcal{A}_2, \Upsilon)\mathcal{A}_3 \\ + S(\mathcal{A}_2, \Upsilon)R(\mathcal{U}, \varsigma)\mathcal{A}_3 - S(\varsigma, \Upsilon)R(\mathcal{U}, \mathcal{A}_2)\mathcal{A}_3 + S(\mathcal{A}_2, \mathcal{A}_3)R(\mathcal{U}, \Upsilon)\varsigma \\ - S(\varsigma, \mathcal{A}_3)R(\mathcal{U}, \Upsilon)\mathcal{A}_2 = 0. \end{aligned} \quad (5.3)$$

Taking inner product with ς , using (3.3) and putting $\Upsilon = \mathcal{A}_3 = \varsigma$, in (5.3), we obtain

$$\begin{aligned} \left[\frac{qr - 2\lambda_1}{2\rho} \right] [\mathfrak{g}(\mathcal{A}_2, R(\mathcal{U}, \varsigma)\varsigma) - \eta(\mathcal{A}_2)\mathfrak{g}(\varsigma, R(\mathcal{U}, \varsigma)\varsigma) - \eta(\mathcal{U})\eta(R(\mathcal{A}_2, \varsigma)\varsigma) \\ + \eta(\mathcal{A}_2)\eta(R(\mathcal{U}, \varsigma)\varsigma) - \eta(R(\mathcal{U}, \mathcal{A}_2)\varsigma) + \eta(\mathcal{A}_2)\eta(R(\mathcal{U}, \varsigma)\varsigma) - \eta(R(\mathcal{U}, \varsigma)\mathcal{A}_2))] \\ + \frac{\lambda_2}{\rho} \{ \eta(\mathcal{U})\eta(R(\mathcal{A}_2, \varsigma)\varsigma) - \eta(\mathcal{A}_2)\eta(R(\mathcal{U}, \varsigma)\varsigma) + \eta(R(\mathcal{U}, \mathcal{A}_2)\varsigma) - \eta(\mathcal{A}_2)\eta(R(\mathcal{U}, \varsigma)\varsigma) \\ + \eta(R(\mathcal{U}, \varsigma)\mathcal{A}_2) \} = 0. \end{aligned} \quad (5.4)$$

Using (2.3) and (2.4) in (5.4) we obtain

$$\begin{aligned} \left[\frac{2\lambda_1 - qr + \lambda_2}{\rho} \right] \mathfrak{g}(\mathcal{A}_2, \mathcal{U}) - \eta(\mathcal{U})\eta(\mathcal{A}_2) = 0, \\ \lambda_1 = \frac{qr - \lambda_2}{2}. \end{aligned}$$

Then by (3.3), we have

$$S(\mathcal{A}_1, \mathcal{A}_2) = \frac{\lambda_2}{\rho} \left[\frac{1}{2} \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_2) - \eta(\mathcal{A}_1)\eta(\mathcal{A}_2) \right]. \quad (5.5)$$

Substituting $\mathcal{A}_2 = \varsigma$ in (5.5), we obtain

$$S(\mathcal{A}_1, \varsigma) = - \left[\frac{\lambda_2}{\rho} \right] \eta(\mathcal{A}_1). \quad (5.6)$$

Comparing this with (2.5) results

$$\lambda_2 = \rho(1 - n), \quad \lambda_1 = \frac{(qr - \rho(1 - n))}{2}.$$

Hence, we have

Theorem 5.1 *If a nearly-Sasakian η -RYS satisfies the condition $S \cdot R = 0$, then $\lambda_2 = \rho(1 - n)$, $\lambda_1 = \frac{qr - \rho(1 - n)}{2}$ and the underlying manifold is η -Einstein.*

6. η -RYS ON EINSTEIN-SEMISYMMETRIC NSMs

Consider an Einstein-semisymmetric NSM Ω . By Definition 2.2,

$$E(R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3, \mathcal{W}) + E(\mathcal{A}_3, R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{W}) = 0. \quad (6.1)$$

Substituting $\mathcal{A}_1 = \varsigma$ and using (2.8), we get

$$S(R(\varsigma, \mathcal{A}_2)\mathcal{A}_3, \mathcal{W}) + S(\mathcal{A}_3, R(\varsigma, \mathcal{A}_2)\mathcal{W}) - \left(\frac{r}{n}\right)[\mathfrak{g}(R(\varsigma, \mathcal{A}_2)\mathcal{A}_3, \mathcal{W}) + \mathfrak{g}(\mathcal{A}_3, R(\varsigma, \mathcal{A}_2)\mathcal{W})] = 0. \quad (6.2)$$

If the Einstein-semisymmetric NSM admits an η -RYS, then by (3.3) we get

$$\begin{aligned} & \left[\frac{qr - 2\lambda_1}{2\rho}\right] \mathfrak{g}(R(\varsigma, \mathcal{A}_2)\mathcal{A}_3, \mathcal{W}) - \frac{\lambda_2}{\rho} \eta(R(\varsigma, \mathcal{A}_2)\mathcal{A}_3) \eta(\mathcal{W}) - \left[\frac{r}{n}\right] \mathfrak{g}(R(\varsigma, \mathcal{A}_2)\mathcal{A}_3, \mathcal{W}) \\ & + \left[\frac{qr - 2\lambda_1}{2\rho}\right] \mathfrak{g}(\mathcal{A}_3, R(\varsigma, \mathcal{A}_2)\mathcal{W}) - \frac{\lambda_2}{\rho} \eta(R(\varsigma, \mathcal{A}_2)\mathcal{W}) \eta(\mathcal{A}_3) \\ & - \frac{r}{n} \mathfrak{g}(\mathcal{A}_3, R(\varsigma, \mathcal{A}_2)\mathcal{W}) = 0. \end{aligned} \quad (6.3)$$

Now by (2.4) and taking $\mathcal{W} = \varsigma$, we obtain

$$-\frac{\lambda_2}{\rho} [\mathfrak{g}(\mathcal{A}_2, \mathcal{A}_3) \eta(\mathcal{W}) - 2\eta(\mathcal{A}_2) \eta(\mathcal{A}_3) \eta(\mathcal{W}) + \eta(\mathcal{A}_3) \mathfrak{g}(\mathcal{A}_2, \mathcal{W})] = 0, \quad (6.4)$$

which implies $\lambda_2 = 0$. Then, by theorem (3.1), we have $\lambda_1 = \rho(1 - n) + \frac{qr}{2}$. Thus, we have

Theorem 6.1 *A nearly-Sasakian η -RYS cannot be Einstein-semisymmetric and $\lambda_1 = \rho(1 - n) + \frac{qr}{2}$.*

7. η -RYS ON 3-DIMENSIONAL NSM SATISFYING $Q \cdot R = 0$

Let Ω be a 3-dimensional NSM admitting η -RYS that satisfies $Q \cdot R = 0$. Then

$$Q(R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3 - R(Q\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3 - R(\mathcal{A}_1, Q\mathcal{A}_2)\mathcal{A}_3 - R(\mathcal{A}_1, \mathcal{A}_2)Q\mathcal{A}_3) = 0, \quad (7.1)$$

for all $\mathcal{A}_1, \mathcal{A}_2, \mathcal{A}_3$ on Ω . In view of (3.4), Eq. (7.1) takes the form

$$\begin{aligned} & \frac{\lambda_2}{\rho} [\eta(\mathcal{A}_1) R(\varsigma, \mathcal{A}_2)\mathcal{A}_3 - \eta(R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3) \varsigma + \eta(\mathcal{A}_2) R(\mathcal{A}_1, \varsigma)\mathcal{A}_3 + \eta(\mathcal{A}_3) R(\mathcal{A}_1, \mathcal{A}_2)\varsigma] \\ & - \frac{(qr - 2\lambda_1)}{2\rho} [2R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3] = 0. \end{aligned}$$

Taking inner product with ς and using the equations (2.3) and (2.4), above equation reduces to

$$\eta(R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3) \left[\frac{2\lambda_1 - \lambda_2 - qr}{\rho} \right] + \frac{\lambda_2}{\rho} [\eta(\mathcal{A}_1) \mathfrak{g}(\mathcal{A}_2, \mathcal{A}_3) - \eta(\mathcal{A}_2) \mathfrak{g}(\mathcal{A}_1, \mathcal{A}_3)] = 0. \quad (7.2)$$

Taking $\mathcal{A}_1 = \varsigma$ in (7.2), we get

$$\frac{(2\lambda_1 - qr)}{\rho} [\mathfrak{g}(\mathcal{A}_2, \mathcal{A}_3) - \eta(\mathcal{A}_2) \eta(\mathcal{A}_3)] = 0.$$

Upon simplification above equation yields

$$\lambda_1 = \frac{qr}{2}.$$

Thus, we have:

Theorem 7.1 *If a 3-dimensional NSM admitting η -RYS satisfies $Q \cdot R = 0$, then the Ricci soliton is*

1. *expanding, for $\frac{qr}{2} > 0$,*
2. *shrinking, for $\frac{qr}{2} < 0$.*

8. η -RYS ON 3-DIMENSIONAL NSM SATISFYING $R(\mathcal{A}, \varsigma) \cdot P - P(\mathcal{A}, \varsigma) \cdot R = 0$

Let M be a 3-dimensional NSM admitting η -RYS that satisfies

$$R(\mathcal{A}, \varsigma) \cdot P - P(\mathcal{A}, \varsigma) \cdot R = 0 \quad (8.1)$$

We require the following definition:

Definition 8.1 *The projective curvature tensor P in a 3-dimensional NSM is defined by*

$$P(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3 = R(\mathcal{A}_1, \mathcal{A}_2)\mathcal{A}_3 - \frac{1}{n}[S(\mathcal{A}_2, \mathcal{A}_3)\mathcal{A}_1 - S(\mathcal{A}_1, \mathcal{A}_3)\mathcal{A}_2]. \quad (8.2)$$

The following results are due to (7.1):

$$P(\mathcal{A}_1, \varsigma) = \left(\frac{2n\rho - qr + 2\lambda_1 + 2\lambda_2}{2n\rho} \right) \eta(\mathcal{A}_2)\mathcal{A}_1 - g(\mathcal{A}_1, \mathcal{A}_2)\varsigma + \frac{1}{n}S(\mathcal{A}_1, \mathcal{A}_2)\varsigma, \quad (8.3)$$

$$P(\mathcal{A}_1, \mathcal{A}_2)\varsigma = \left(\frac{2n\rho - qr + 2\lambda_1 + 2\lambda_2}{2n\rho} \right) (\eta(\mathcal{A}_2)\mathcal{A}_1 - \eta(\mathcal{A}_1)\mathcal{A}_2), \quad (8.4)$$

$$P(\mathcal{A}_1, \varsigma)\varsigma = \left(\frac{2n\rho - qr + 2\lambda_1 + 2\lambda_2}{2n\rho} \right) \mathcal{A}_1 - \eta(\mathcal{A}_1)\varsigma. \quad (8.5)$$

From (8.1), we have

$$\begin{aligned} & R(\mathcal{A}, \varsigma)P(\mathcal{U}, \Upsilon)\mathcal{W} - P(R(\mathcal{A}, \varsigma)\mathcal{U}, \Upsilon)\mathcal{W} - P(\mathcal{U}, R(\mathcal{A}, \varsigma)\Upsilon)\mathcal{W} \\ & - P(\mathcal{U}, \Upsilon)R(\mathcal{A}, \varsigma)\mathcal{W} - P(\mathcal{A}, \varsigma)R(\mathcal{U}, \Upsilon)\mathcal{W} + R(P(\mathcal{A}, \varsigma)\mathcal{U}, \Upsilon)\mathcal{W} \\ & + R(\mathcal{U}, P(\mathcal{A}, \varsigma)\Upsilon)\mathcal{W} + R(\mathcal{U}, \Upsilon)P(\mathcal{A}, \varsigma)\mathcal{W} = 0. \end{aligned}$$

Substituting $\mathcal{U} = \mathcal{W} = \varsigma$, yields

$$\begin{aligned} & R(\mathcal{A}, \varsigma)P(\varsigma, \Upsilon)\varsigma - P(R(\mathcal{A}, \varsigma)\varsigma, \Upsilon)\varsigma - P(\varsigma, R(\mathcal{A}, \varsigma)\Upsilon)\varsigma - P(\varsigma, \Upsilon)R(\mathcal{A}, \varsigma)\varsigma \\ & - P(\mathcal{A}, \varsigma)R(\varsigma, \Upsilon)\varsigma + R(P(\mathcal{A}, \varsigma)\varsigma, \Upsilon)\varsigma \\ & + R(\varsigma, P(\mathcal{A}, \varsigma)\Upsilon)\varsigma + R(\varsigma, \Upsilon)P(\mathcal{A}, \varsigma)\varsigma = 0. \end{aligned} \quad (8.6)$$

Now, using (2.3), (2.4), (8.3)-(8.5) in (8.6), we obtain

$$\frac{2}{n}S(\Upsilon, \mathcal{A})\varsigma - g(\Upsilon, \mathcal{A})\varsigma + \eta(\mathcal{A})\eta(\Upsilon)\varsigma \left[\frac{n\rho - qr + 2\lambda_1 + 2\lambda_2}{n\rho} \right] = 0.$$

Applying inner product with ς in the aforementioned equation, we achieve

$$S(\Upsilon, \mathcal{A}) = \left(\frac{n}{2} \right) g(\Upsilon, \mathcal{A}) - \eta(\mathcal{A})\eta(\Upsilon) \left(\frac{n\rho - qr + 2\lambda_1 + 2\lambda_2}{2\rho} \right). \quad (8.7)$$

Putting $\mathcal{A} = \varsigma$ in (8.7) and using (2.5), we get $\lambda_1 = (1 - n)\rho + \frac{qr}{2} - \lambda_2$.

Hence (8.7) reduces to

$$S(\Upsilon, \mathcal{A}) = \left(\frac{n}{2} \right) \mathfrak{g}(\Upsilon, \mathcal{A}) + \left(\frac{n}{2} - 1 \right) \eta(\mathcal{A})\eta(\Upsilon).$$

Hence, we have

Theorem 8.1 *If a NSM of dimension 3 admitting η -RYS satisfies the relation*

$$R(\mathcal{A}, \varsigma) \cdot P - P(\mathcal{A}, \varsigma) \cdot R = 0,$$

then the Ricci soliton is steady and the manifold is Einstein.

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