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# Certain Subclasses of Bi-Univalent Functions Defined by q-Analogue of Ruscheweyh Differential Operator

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ABSTRACT: In this paper, we find a new subclasses of the function class  $\sum$  of bi-univalent functions defined in the open unit disk, which are associated with the q-analogue of Ruscheweyh differential operator and satisfy some subordination conditions. Furthermore, we find estimates on the Taylor-Maclaurin coefficients  $|v_2|$  and  $|v_3|$  for functions in the new subclasses introduced here.

Key Words: Analytic functions, univalent functions, bi-univalent, Starlike and convex functions, q-Ruscheweyh differential operator.

### Contents

#### 1. Introduction

Let  $\mathcal{A}$  denote the class of analytic functions of the form

$$\psi(u) = u + \sum_{j=2}^{\infty} v_j u^j \tag{1.1}$$

normalized by the conditions  $\psi(0) = 0 = \psi'(0) - 1$ , defined in the open unit disk

$$U = \{ u \in \mathcal{C} : |u| < 1 \}.$$

Let  $\mathcal{M}$  be the subclass of  $\mathcal{A}$  consisting of function of the form (1) which are also univalent in U. Consider an analytic function  $\xi$  with positive real part in the unit disk U,  $\xi(0) = 1, \xi'(0) > 0$  and  $\xi$  maps U onto a region starlike with respect to 1 and symmetric with respect to the real axis. In the sequel, it is assumed that such a function has a series expansion of the form

$$\zeta(u) = 1 + B_1 u + B_2 u^2 + B_3 u^3 + \dots, (B_1 > 0). \tag{1.2}$$

In particular, for the class of strongly starlike functions of order  $\alpha(0 < \alpha \le 1)$ , the function  $\zeta$  is given by

$$\zeta(u) = \left[ \frac{1+u}{1-u} \right]^{\alpha} = 1 + 2\alpha u + 2\alpha^2 u^2 + \dots \qquad (0 < \alpha \le 1), \tag{1.3}$$

which gives  $B_1 = 2\alpha$  and  $B_2 = 2\alpha^2$  and on the other hand, for the class of starlike functions of order  $\beta(0 \le \beta < 1)$ ,

$$\zeta(u) = \frac{1 + (1 - 2\beta)u}{1 - u} = 1 + 2(1 - \beta)u + 2(1 - \beta)u^2 + \dots \qquad (0 \le \beta < 1), \tag{1.4}$$

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we have  $B_1 = B_2 = 2(1 - \beta)$ .

A function  $\psi \in \mathcal{A}$  is said to be bi-univalent in U if both  $\psi$  and  $\psi^{-1}$  are univalent in U. Let  $\Sigma$  denote the class of bi-univalent functions defined in the unit disk U. Since  $\psi \in \Sigma$  has the Maclaurian series given by (1), a computation shows that its inverse  $\phi = \psi^{-1}$  has the expansion

$$\phi(w) = \psi^{-1}(w) = w - a_2 w^2 + (2a_2^2 - a_3)w^3 + \cdots$$
(1.5)

Several authors have introduced and investigated subclasses of bi-univalent functions and obtained bounds for the initial coefficients (See [2,3,4,8,11,12,13]).

For 0 < q < 1, the Jackson's q-derivative of a function  $\psi(u) \in \mathcal{A}$  is given by ([5]):

$$D_{q}\psi(u) = \begin{cases} \frac{\psi(u) - \psi(qu)}{(1-q)u} & for \quad u \neq 0; \\ \psi'(0) & for \quad u = 0. \end{cases}$$
 (1.6)

For  $\psi(u)$  of the form (1), we have

$$D_q \psi(u) = 1 + \sum_{j=2}^{\infty} [j]_q v_j u^{j-1}, \qquad (1.7)$$

where

$$[j]_q = \frac{1 - q^j}{1 - q} \quad (0 < q < 1; j \in \mathbb{N} = \{1, 2, \dots\}). \tag{1.8}$$

Kanas and Raducanu [7] (See also Aldweby and Darus [1]) defined the q-analogue of Ruscheweyh operator by:

$$R_q^{\lambda} \psi(u) = u + \sum_{j=2}^{\infty} \frac{[j+\lambda-1]_q!}{[\lambda]_q![j-1]_q!} v_j u^j \quad (0 < q < 1; \lambda \ge 0),$$
(1.9)

where

$$[j]_q! = \begin{cases} [j]_q [j-1]_q \dots [1]_q, & j \in \mathbb{N}; \\ 1, & j = 0. \end{cases}$$
 (1.10)

From (9), we obtain that

$$R_q^0\psi(u) = \psi(u)$$
 and  $R_q^1\psi(u) = uD_q\psi(u)$ ,

and

$$\lim_{q \to 1^{-}} R_{q}^{\lambda} \psi(u) = u + \sum_{j=2}^{\infty} \frac{[j+\lambda-1]_{q}!}{[\lambda]_{q}![j-1]_{q}!} v_{j} u^{j} = R^{\lambda} \psi(u), \tag{1.11}$$

where  $R^{\lambda}$  is the Ruscheweyh differential operator [10].

# 2. Bi-Univalent Function Class $\mathcal{N}\Sigma_a^{\lambda}(\delta,\zeta)$

In this section, we introduce a subclass  $\mathcal{N}\Sigma_q^{\lambda}(\delta,\zeta)$  of  $\Sigma$  and find the estimate on the coefficients  $|v_2|$  and  $|v_3|$  for the functions in this new subclass, by subordination. Throughout our study, unless otherwise stated, we let

$$0 < \delta < 1$$
 and  $0 < q < 1$ .

**Definition 2.1** For  $0 \le \delta \le 1$ , a function  $\psi \in \Sigma$  of the form (1.1) is said to be in the class  $\mathcal{N}\Sigma_q^{\lambda}(\delta,\zeta)$ , if the following subordination hold:

$$(1 - \delta) \frac{u D_q R_q^{\lambda} \psi(u)}{R_q^{\lambda} \psi(u)} + \delta \frac{D_q(u D_q R_q^{\lambda} \psi(u))}{D_q(R_q^{\lambda} \psi(u))} \prec \zeta(u)$$
(2.1)

and

$$(1 - \delta) \frac{w D_q R_q^{\lambda} \phi(w)}{R_q^{\lambda} \phi(w)} + \delta \frac{D_q(w D_q R_q^{\lambda} \phi(w))}{D_q(R_q^{\lambda} \phi(w))} \prec \zeta(w), \tag{2.2}$$

where  $u, w \in U$  and  $\phi$  is given by (5).

Note that if  $\lambda=0$  and  $q\to 1^-$  the class  $\mathcal{N}\Sigma_q^{\lambda}(\delta,\zeta)$  reduces to class  $M_{\Sigma}(\alpha,\lambda),\ 0<\alpha\leq 1$  and  $\lambda\geq 0$  studied by Xiao-Fei Li and Au-Ping Wang [14].

If  $\lambda = 1$  and  $q \to 1^-$  the class  $\mathcal{N}\Sigma_q^{\lambda}(\delta, \zeta)$  satisfying the subordination  $\frac{u\psi'(u)}{\psi(u)} \prec \zeta(u)$  and  $1 + \frac{u\psi''^{(u)}}{\psi'(u)} \prec \zeta(u)$  studied by the class of Ma and Minda [9] starlike and convex function respectively.

If  $\lambda = 1$  and  $q \to 1^-$  the class  $\mathcal{N}\Sigma_q^{\lambda}(\delta, \zeta)$  reduces to class  $M_{\sigma}(\alpha, \varphi)$  studied by Jothi Latha and Cynthiya Margaret Indrani [6].

**Lemma 2.1** If a function  $p \in \mathcal{P}$  is given by

$$p(u) = 1 + p_1 u + p_2 u^2 + \cdots \quad (u \in U),$$

then

$$|p_i| \le 2 \quad (i \in \mathbb{N}),$$

where is  $\mathcal{P}$  is the family of all functions p, analytic in U, for which

$$p(0) = 1$$
 and  $\Re(p(u)) > 0$   $(u \in U)$ .

**Theorem 2.1** If  $\psi$  given by (1) is in the class  $\mathcal{N}\Sigma_q^{\lambda}(\delta,\zeta)$ , then

$$|v_{2}| \leq \frac{B_{1}\sqrt{B_{1}}}{\left|\left(\left[1+\delta([3]_{q})-1\right]\frac{(\lambda+1)_{q}(\lambda+2)_{q}}{[2]_{q}!}([3]_{q}-1)-\left[1+\delta([2]_{q}^{2}-1)\right](\lambda+1)_{q}^{2}\right|} (2.3)$$

$$\sqrt{([2]_{q}-1))B_{1}^{2}+\left[1+\delta([2]_{q})-1\right](B_{1}-B_{2})(\lambda+1)_{q}^{2}([2]_{q}-1)^{2}}$$

and

$$|v_3| \le \frac{B_1}{[1 + \delta([3]_q) - 1] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} ([3]_q - 1)} + \left(\frac{B_1}{[1 + \delta([2]_q) - 1](\lambda + 1)_q([2]_q - 1)}\right)^2, \tag{2.4}$$

where  $0 \le \delta \le 1$ .

**Proof:** Let  $\psi \in \mathcal{N}\Sigma_q^{\lambda}(\delta,\zeta)$  and  $\phi \in \psi^{-1}$ . Then there are analytic functions  $a,b: U \to U$ , with a(0) = 0 = b(0), satisfying

$$(1 - \delta) \frac{u D_q R_q^{\lambda} \psi(u)}{R_q^{\lambda} \psi(u)} + \delta \frac{D_q (u D_q R_q^{\lambda} \psi(u))}{D_q (R_q^{\lambda} \psi(u))} = \zeta(a(u))$$

$$(2.5)$$

and

$$(1 - \delta) \frac{w D_q R_q^{\lambda} \phi(w)}{R_q^{\lambda} \phi(w)} + \delta \frac{D_q(w D_q R_q^{\lambda} \phi(w))}{D_q(R_q^{\lambda} \phi(w))} = \zeta(b(w)). \tag{2.6}$$

Define the functions p(u) and q(u) by

$$p(u) := \frac{1 + a(u)}{1 - a(u)} = 1 + p_1 u + p_2 u^2 + \cdots$$

and

$$q(u) := \frac{1 + b(u)}{1 - b(u)} = 1 + q_1 u + q_2 u^2 + \cdots$$

or, equivalently,

$$a(u) := \frac{p(u) - 1}{p(u) + 1} = \frac{1}{2} \left[ p_1 u + \left( p_2 - \frac{p_1^2}{2} \right) u^2 + \cdots \right]$$
 (2.7)

and

$$b(u) := \frac{q(u) - 1}{q(u) + 1} = \frac{1}{2} \left[ q_1 u + \left( q_2 - \frac{q_1^2}{2} \right) u^2 + \dots \right]. \tag{2.8}$$

Then p(u) and q(u) are analytic in U with p(0) = 1 = q(0). Since  $a, b : U \to U$ , the functions p(u) and q(u) have a positive real part in U,  $|p_i| \le 2$  and  $|q_i| \le 2$ .

Using (18) and (19) in (16) and (17) respectively, we have

$$(1 - \delta) \frac{u D_q R_q^{\lambda} \psi(u)}{R_q^{\lambda} \psi(u)} + \delta \frac{D_q(u D_q R_q^{\lambda} \psi(u))}{D_q(R_q^{\lambda} \psi(u))} = \zeta \left( \frac{1}{2} \left[ p_1 u + \left( p_2 - \frac{p_1^2}{2} \right) u^2 + \cdots \right] \right)$$
 (2.9)

and

$$(1 - \delta) \frac{w D_q R_q^{\lambda} \phi(w)}{R_q^{\lambda} \phi(w)} + \delta \frac{D_q(w D_q R_q^{\lambda} \phi(w))}{D_q(R_q^{\lambda} \phi(w))} = \zeta \left( \frac{1}{2} \left[ q_1 w + \left( q_2 - \frac{q_1^2}{2} \right) w^2 + \cdots \right] \right). \tag{2.10}$$

In light of (1)-(5), and from (20) and (21), we have

$$1 + [1 + \delta([2]_q - 1)](\lambda + 1)_q([2]_q - 1)v_2u + \left\{ \left( [1 + \delta([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} ([3]_q - 1)v_3 \right) - [1 + \delta([2]_q^2 - 1)](\lambda + 1)_q^2([2]_q - 1)v_2^2 \right\} u^2 + \cdots$$

$$= 1 + \frac{1}{2}B_1p_1u + \left[ \frac{1}{2}B_1(p_2 - \frac{p_1^2}{2}) + \frac{1}{4}B_2p_1^2 \right] u^2 + \cdots$$

and

$$\begin{split} 1 - [1 + \delta([2]_q - 1)](\lambda + 1)_q([2]_q - 1)v_2w + & \left\{ \left( 2[1 + \delta([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} ([3]_q - 1) - [1 + \delta([2]_q^2 - 1)](\lambda + 1)_q^2([2]_q - 1) \right) v_2^2 - [1 + \delta([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} ([3]_q - 1)v_3 \right\} w^2 + \cdots \\ & = 1 + \frac{1}{2} B_1 q_1 w + \left[ \frac{1}{2} B_1 (q_2 - \frac{q_1^2}{2}) + \frac{1}{4} B_2 q_1^2 \right] w^2 + \cdots, \end{split}$$

which yields the following relations:

$$[1 + \delta([2]_q - 1)](\lambda + 1)_q([2]_q - 1)v_2 = \frac{1}{2}B_1p_1$$
(2.11)

$$-\left[1+\delta([2]_q^2-1)\right](\lambda+1)_q^2([2]_q-1)v_2^2 + \left[1+\delta([3]_q-1)\right]\frac{(\lambda+1)_q(\lambda+2)_q}{[2]_q!}([3]_q-1)v_3$$

$$=\frac{1}{2}B_1(p_2-\frac{p_1^2}{2}) + \frac{1}{4}B_2p_1^2$$
(2.12)

$$-[1+\delta([2]_q-1)](\lambda+1)_q([2]_q-1)v_2 = \frac{1}{2}B_1q_1$$
 (2.13)

$$\left(2[1+\delta([3]_q-1)]\frac{(\lambda+1)_q(\lambda+2)_q}{[2]_q!}([3]_q-1)-[1+\delta([2]_q^2-1)](\lambda+1)_q^2([2]_q-1)\right)v_2^2- \\
[1+\delta([3]_q-1)]\frac{(\lambda+1)_q(\lambda+2)_q}{[2]_q!}([3]_q-1)v_3 = \frac{1}{2}B_1(q_2-\frac{q_1^2}{2})+\frac{1}{4}B_2q_1^2.$$
(2.14)

From (22) and (24), it follows that

$$p_1 = -q_1 (2.15)$$

and

$$8[1 + \delta([2]_q - 1)]^2(\lambda + 1)_q^2([2]_q - 1)^2v_2^2 = B_1^2(p_1^2 + q_1^2).$$
(2.16)

From (23), (25) and (27), we obtain

$$v_{2}^{2} = \frac{B_{1}^{3}(p_{2} + q_{2})}{4\left\{\left(\left[1 + \delta([3]_{q} - 1)\right]\frac{(\lambda + 1)_{q}(\lambda + 2)_{q}}{[2]_{q}!}([3]_{q} - 1) - \left[1 + \delta([2]_{q}^{2} - 1)\right](\lambda + 1)_{q}^{2}\right\}}$$

$$([2]_{q} - 1)\right\}B_{1}^{2} + \left[1 + \delta([2]_{q} - 1)\right](B_{1} - B_{2})(\lambda + 1)_{q}^{2}([2]_{q} - 1)^{2}\right\}.$$

Applying Lemma 2.1 to the coefficients  $p_2$  and  $q_2$ , we have

$$|v_{2}| \leq \frac{B_{1}\sqrt{B_{1}}}{\left|\left(\left[1+\delta([3]_{q})-1\right]\frac{(\lambda+1)_{q}(\lambda+2)_{q}}{\left[2\right]_{q}!}\left(\left[3\right]_{q}-1\right)-\left[1+\delta([2]_{q}^{2}-1)\right](\lambda+1)_{q}^{2}\right|} \sqrt{\left([2]_{q}-1\right)B_{1}^{2}+\left[1+\delta([2]_{q})-1\right](B_{1}-B_{2})(\lambda+1)_{q}^{2}([2]_{q}-1)^{2}\right|}.$$
(2.18)

By substracting (25) from (23) and using (26) and (27), we get

$$v_{3} = \frac{B_{1}^{2}(p_{1}^{2} + q_{1}^{2})}{8[1 + \delta([2]_{q}) - 1]^{2}(\lambda + 1)_{q}^{2}([2]_{q} - 1)^{2}} + \frac{B_{1}(p_{2} - q_{2})}{4[1 + \delta([3]_{q} - 1)]\frac{(\lambda + 1)_{q}(\lambda + 2)_{q}}{[2]_{q}!}([3]_{q} - 1)}.$$
(2.19)

Applying Lemma 2.1 once again to the coefficients  $p_1, p_2, q_1$  and  $q_2$ , we get

$$|v_{3}| \leq \frac{B_{1}}{[1 + \delta([3]_{q}) - 1] \frac{(\lambda + 1)_{q}(\lambda + 2)_{q}}{[2]_{q}!} ([3]_{q} - 1)} + \left(\frac{B_{1}}{[1 + \delta([2]_{q}) - 1](\lambda + 1)_{q}([2]_{q} - 1)}\right)^{2}.$$
(2.20)

3. Bi-Univalent Function Class  $\mathcal{F}\Sigma_q^{\lambda}(\mu,\zeta)$ 

**Definition 3.1** For  $0 \le \mu \le 1$ , a function  $\psi \in \Sigma$  of the form (1) is said to be in the class  $\mathcal{F}\Sigma_q^{\lambda}(\mu,\zeta)$ , if the following subordination hold:

$$(1-\mu)\frac{R_q^{\lambda}\psi(u)}{u} + \mu D_q R_q^{\lambda}\psi(u) \prec \zeta(u)$$
(3.1)

and

$$(1-\mu)\frac{R_q^{\lambda}\phi(w)}{w} + \mu D_q R_q^{\lambda}\phi(w) \prec \zeta(w), \tag{3.2}$$

where  $u, w \in U, \phi$  is given by (5) and  $R_q^{\lambda}\psi(u)$  is given by (9).

**Theorem 3.1** Let  $\psi$  given by (1) be in the class  $\mathcal{F}\Sigma_a^{\lambda}(\mu,\zeta)$ . Then

$$|v_2| \le \frac{B_1 \sqrt{B_1}}{\sqrt{\left| \left[1 + \mu([3]_q - 1)\right] \frac{(\lambda + 1)_q (\lambda + 2)_q}{[2]_q!} B_1^2 + \left[1 + \mu([2]_q - 1)\right]^2 (\lambda + 1)_q^2 (B_1 - B_2)\right|}}$$
(3.3)

and

$$|v_3| \le \frac{B_1}{[1 + \mu([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!}} + \left(\frac{B_1}{[1 + \mu([2]_q - 1)](\lambda + 1)_q}\right)^2.$$
(3.4)

**Proof:** Proceeding as in the proof of Theorem 2.1, we can arrive the following relations:

$$[1 + \mu([2]_q - 1)](\lambda + 1)_q v_2 = \frac{1}{2} B_1 p_1, \tag{3.5}$$

$$[1 + \mu([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} v_3 = \frac{1}{2} B_1 \left( p_2 - \frac{p_1^2}{2} \right) + \frac{1}{4} B_2 p_1^2, \tag{3.6}$$

$$-[1 + \mu([2]_q - 1)](\lambda + 1)_q v_2 = \frac{1}{2} B_1 q_1, \tag{3.7}$$

$$2[1 + \mu([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} v_2^2 - [1 + \mu([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} v_3$$

$$= \frac{1}{2} B_1 \left( q_2 - \frac{q_1^2}{2} \right) + \frac{1}{4} B_2 q_1^2$$
(3.8)

From (36) and (38) it follows that

$$p_1 = -q_1 \tag{3.9}$$

and

$$8[1 + \mu([2]_q - 1)]^2 (\lambda + 1)_q^2 v_2^2 = B_1^2 (p_1^2 + q_1^2).$$
(3.10)

From (37), (39) and (41), we obtain

$$v_2^2 = \frac{B_1^3(p_2 + q_2)}{4\left\{ [1 + \mu([3]_q - 1)] \frac{(\lambda + 1)_q(\lambda + 2)_q}{[2]_q!} B_1^2 + (B_1 - B_2)[1 + \mu([2]_q - 1)]^2(\lambda + 1)_q^2 \right\}}.$$
 (3.11)

Applying Lemma 2.1 to the coefficient  $p_2$  and  $q_2$ , we immediately get the desired estimate on  $|v_2|$  as asserted in (34). By subtracting (39) from (37) and using (40) and (41), we get

$$v_3 = \frac{B_1(p_2 - q_2)}{4[1 + \mu([3]_q - 1)] \frac{(\lambda + 1)(\lambda + 2)_q}{[2]_q!}} + \frac{B_1^2(p_1^2 + q_1^2)}{8[1 + \mu([2]_q - 1)]^2(\lambda + 1)_q^2}.$$
 (3.12)

Applying Lemma 2.1 to the coefficients  $p_1, p_2, q_1$  and  $q_2$ , we get the desired estimate on  $|v_3|$  as asserted in (35).

# 4. Conclusion

We considered the q-Analogue of Ruscheweyh differential operator and defined a new subclasses of the bi-univalent functions in open unit disk. We investigated Taylor-Maclaurin coefficients  $|v_2|$  and  $|v_3|$  for functions belonging to this new subclasses and its subclasses and discussed some geometric properties of these subclasses.

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