



Study of some unified integrals associated with extended k-generalized Mittag-Leffler function

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ABSTRACT: Many authors established a large range of integral formulas including various special functions throughout the years. We evolve certain integrals formulas involving extended k -generalized Mittag-Leffler function using the extension of beta function in this paper. Some special cases also discussed for our main results..

Key Words: Extended Mittag-Leffler function, extended k -Mittag-Leffler function, extended k -Beta function, extended k -Gamma function, k -Pochhammer symbol.

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1. Introduction

During the past few decades special function and fractional calculus become active research area being diverse applications in physics and engineering. Special functions play a important role to solve both integral and differential equations of fractional order and Mittag-Leffler function is one of the tool which is used frequently by authors [4,10,11,16,26,14,15,16,17,27,19,20,21,22,23,24,25,26,27]. It has enlarge applications in the field of Applied Sciences, Biology and many other areas.

The elementary Mittag-Leffler function is defined and represented as

$$E_r(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(nr+1)}, (r > 0, z \in \mathbb{C}) \quad (1.1)$$

let $r = 2$ and $r = 4$ then equation (1.1) become

$$E_2(z) = \cosh(\sqrt{z}), (z \in \mathbb{C}) \quad (1.2)$$

$$E_4(z) = \frac{1}{2} \left[\cos(z^{\frac{1}{4}}) + \cosh(z^{\frac{1}{4}}) \right], (z \in \mathbb{C}) \quad (1.3)$$

Its generalization was given by Wiman [1]

$$E_{r,s}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(nr+s)}, (r, s, z \in \mathbb{C}, \Re(r) > 0, \Re(s) > 0) \quad (1.4)$$

After Wiman, Prabhakar [2] introduced generalization of Mittag-Leffler function as

$$E_{r,s}^{\rho}(z) = \sum_{n=0}^{\infty} \frac{(\rho)_n}{\Gamma(nr+s)} \frac{z^n}{n!}, (r, s, z, \rho \in \mathbb{C}; \Re(r) > 0, \Re(s) > 0) \quad (1.5)$$

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In 2007 Shukla and Prajapati [3] extended the definition of Prabhakar function as

$$E_{r,s}^{\rho,\sigma}(z) = \sum_{n=0}^{\infty} \frac{(\rho)_{n\sigma}}{\Gamma(nr+s)} \frac{z^n}{n!}, (r, s, z, \rho \in \mathbb{C}; \Re(r) > 0, \Re(s) > 0, \Re(\rho) > 0, \sigma \in (0, 1) \cup \Re) \quad (1.6)$$

G.A. Dorrego et.al. [4] in 2012, use the concept of k -Gamma function to define Mittag-Leffler function as

$$E_{k,r,s}^{\rho}(z) = \sum_{n=0}^{\infty} \frac{(\rho)_{n,k}}{\Gamma_k(nr+s)} \frac{z^n}{n!}, (r, s, z, \rho \in \mathbb{C}; k > 0; \Re(r) > 0, \Re(s) > 0, \Re(\rho) > 0) \quad (1.7)$$

where $(\rho)_{n,k}$ and $\Gamma_k(z)$ are well known k -Pochhammer symbol, and k -Gamma function.

The above k -Mittag Leffler function is further extended and generalized by Gehlot et.al. [5,10,21]

$$GE_{k,r,s}^{\rho,\sigma}(z) = \sum_{n=0}^{\infty} \frac{(\rho)_{n\sigma,k}}{\Gamma_k(nr+s)} \frac{z^n}{n!}, (k \in \mathbb{R}; z, r, s, \rho \in \mathbb{C}; \Re(r), \Re(s), \Re(\rho) > 0, \sigma \in (0, 1) \cup \Re) \quad (1.8)$$

In 2018, Rahman et.al. [6] discussed on extension of k -Mittag leffler function as

$$E_{k,r,s}^{\rho;c}(z; p) = \sum_{n=0}^{\infty} \frac{B_k(\rho + nk, c - \rho; p)(c)_{n,k}}{B_k(\rho, c - \rho)\Gamma_k(nr+s)} \frac{z^n}{n!} \quad (1.9)$$

For

$$\frac{(\rho)_{n,k}}{(c)_{n,k}} = \frac{B_k(\rho + nk, c - \rho)}{B_k(\rho, c - \rho)} \quad (1.10)$$

where $B_k(r, s; p)$ is an extension of the k -Beta function given as

$$B_k(r, s) = \frac{1}{k} \int_0^1 u^{\frac{r}{k}-1} (1-u)^{\frac{s}{k}-1} e^{-\frac{p^k}{ku(1-u)}} du. \quad (1.11)$$

Shilpi jain et.al. [7] 2023 define extended k -generalized Mittag-Leffler function represents as follows:

$$E_{k,r,s}^{\rho,\sigma;c}(z; p) = \sum_{n=0}^{\infty} \frac{B_k(\rho + n\sigma k, c - \rho; p)}{B_k(\rho, c - \rho)} \frac{(c)_{n\sigma,k}}{\Gamma_k(nr+s)} \frac{z^n}{n!} \quad (1.12)$$

where $r, s, z, \rho \in \mathbb{C}; k > 0; \Re(r) > 0, \Re(s) > 0, \Re(\rho) > 0$ and $\sigma \in (0, 1) \cup \Re; p \geq 0$.

Fox and Wright [8], describe asymptotic expansion of wright hyper-geometric function such as generalized hypergeometric series ${}_pF_q$ as follows

$${}_p\Psi_q \left[\begin{matrix} (l_1, L_1), \dots, (l_p, L_p); \\ (m_1, M_1), \dots, (m_q, M_q); \end{matrix} z \right] = \sum_{n=0}^{\infty} \frac{\prod_{j=1}^p \Gamma(l_j + L_j n)}{\prod_{j=1}^q \Gamma(m_j + M_j n)} \frac{z^n}{n!} \quad (1.13)$$

where the coefficients $L_1, \dots, L_p$ and $M_1, \dots, M_q$ are positive real numbers such that

(i) $1 + \sum_{j=1}^q M_j - \sum_{j=1}^p L_j > 0$ and $0 < |z| < \infty; z \neq 0$

(ii) $1 + \sum_{j=1}^q M_j - \sum_{j=1}^p L_j = 0$ and $0 < |z| < L_1^{-L_1} \dots L_p^{-L_p} M_1^{-B_1} \dots M_q^{-M_q}$

The Wright function is a special case of (1.13)

$${}_p\Psi_q \left[\begin{matrix} (l_1, 1), \dots, (l_p, 1); \\ (m_1, 1), \dots, (m_q, 1); \end{matrix} z \right] = \frac{\prod_{j=1}^p \Gamma(l_j)}{\prod_{j=1}^q \Gamma(m_j)} {}_pF_q \left[\begin{matrix} l_1, \dots, l_p; \\ m_1, \dots, m_q; \end{matrix} z \right], \quad (1.14)$$

where ${}_pF_q$ is the generalized hypergeometric function defined by

$${}_pF_q \left[\begin{matrix} l_1, \dots, l_p; \\ m_1, \dots, m_q; \end{matrix} z \right] = \sum_{n=0}^{\infty} \frac{(l_1)_n \dots (l_p)_n}{(m_1)_n \dots (m_q)_n} \frac{z^n}{n!} = {}_pF_q(l_1, \dots, l_p; m_1, \dots, m_q; z) \quad (1.15)$$

Lavoie and Trottier integral formula [9,22,25] will be required for our present investigation .

$$\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} d\lambda = \left(\frac{2}{3}\right)^{2\xi} \frac{\Gamma(\xi)\Gamma(\eta)}{\Gamma(\xi+\eta)} \quad (1.16)$$

where $\Re(\xi) > 0$ and $\Re(\eta) > 0$

2. Integral formulas for Extended k -Generalized Mittag-Leffler function

Theorem 1 For $k > 0; z, r, s, \rho \in \mathfrak{C}; \Re(r) > 0, \Re(s) > 0, \Re(\rho) > 0, \Re(\xi) > 0, \Re(\eta) > 0$ and $\sigma \in (0, 1) \cup \Re; p \geq 0$ following integral formula holds true

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma;c} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (\rho, \sigma k), (\xi, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (2.1)$$

Proof: using equ. (1.12) in the left hand side of equ. (2.1)

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma;c} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ = \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n \lambda^n \left(1 - \frac{\lambda}{3}\right)^{2n}}{n!} d\lambda \end{aligned}$$

changing the order of integration and summation

$$= \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \int_0^1 \lambda^{n+\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2n+2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} d\lambda$$

Now using (1.16) in above equation

$$\begin{aligned} &= \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \left(\frac{2}{3}\right)^{2n+2\xi} \frac{\Gamma(n+\xi)\Gamma(\eta)}{\Gamma(n+\xi+\eta)} \\ &= \left(\frac{2}{3}\right)^{2\xi} \Gamma(\eta) \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \frac{\Gamma(n+\xi)}{\Gamma(n+\xi+\eta)} \left(\frac{2}{3}\right)^{2n} \end{aligned}$$

using equ. (1.10) in above

$$= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) \sum_{n=0}^{\infty} \frac{(\rho)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{(\xi)_n}{(\xi+\eta)_n} \frac{(4y/9)^n}{n!}$$

thus we obtain the required result (2.1)

Theorem 2 For $k > 0; z, r, s, \rho \in \mathfrak{C}; \Re(r) > 0, \Re(s) > 0, \Re(\rho) > 0, \Re(\xi) > 0, \Re(\eta) > 0$ and $\sigma \in (0, 1) \cup \Re; p \geq 0$ following integral formula holds true

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma;c} \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (\rho, \sigma k), (\eta, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (2.2)$$

Proof: using equ. (1.12) in left hand side of equ. (2.2)

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma;c} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ = \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n (1-\lambda)^{2n} \left(1 - \frac{\lambda}{4}\right)^n}{n!} d\lambda \end{aligned}$$

changing the order of integration and summation

$$= \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2n+2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{n+\eta-1} d\lambda$$

Now using (1.16) in above equation

$$\begin{aligned} &= \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \left(\frac{2}{3}\right)^{2\xi} \frac{\Gamma(\xi)\Gamma(n+\eta)}{\Gamma(n+\xi+\eta)} \\ &= \left(\frac{2}{3}\right)^{2\xi} \Gamma(\xi) \sum_{n=0}^{\infty} \frac{B_k(\rho+n\sigma k, c-\rho)}{B_k(\rho, c-\rho)} \frac{(c)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{y^n}{n!} \frac{\Gamma(n+\eta)}{\Gamma(n+\xi+\eta)} \end{aligned}$$

using equ. (1.10) in above

$$= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) \sum_{n=0}^{\infty} \frac{(\rho)_{n\sigma, k}}{\Gamma_k(nr+s)} \frac{(\eta)_n}{(\xi+\eta)_n} \frac{(y)^n}{n!}$$

thus we obtain the required result (2.2)

3. Special Cases

Corollary 1 If we put $c = 1$ in equation (2.1) and use equation (1.8), we get

$$\begin{aligned} &\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ &= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (\rho, \sigma k), (\xi, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.1)$$

Corollary 2 If we put $c = \sigma = 1$ in equation (2.1) and use equation (1.7), we get

$$\begin{aligned} &\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ &= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (\rho, k), (\xi, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.2)$$

Corollary 3 If we put $c = \sigma = k = 1$ in equation (2.1) and use equation (1.5), we get

$$\begin{aligned} &\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{r,s}^{\rho} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ &= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (\rho, 1), (\xi, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.3)$$

Corollary 4 If we put $c = \sigma = k = \rho = 1$ in equation (2.1) and use equation (1.4), we get

$$\begin{aligned} &\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{r,s} \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ &= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (1, 1), (\xi, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.4)$$

Corollary 5 If we put $c = \sigma = k = \rho = s = 1$ in equation (2.1) and use equation (1.1), we get

$$\begin{aligned} &\int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_r \left[y\lambda \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ &= \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (1, 1), (\xi, 1); \\ (1, r), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.5)$$

Corollary 6 If we put $c = \sigma = k = \rho = s = 1$ and $r = 2$ in equation (2.1) and use equation (1.2), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \cosh \left[\sqrt{y\lambda} \left(1 - \frac{\lambda}{3}\right)^2 \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (1, 1), (\xi, 1); \\ (1, 2), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.6)$$

Corollary 7 If we put $c = \sigma = k = \rho = s = 1$ and $r = 4$ in equation (2.1) and use equation (1.3), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \left\{ \cos \left(\lambda y \left(1 - \frac{\lambda}{3}\right)^2 \right)^{\frac{1}{4}} + \cosh \left(\lambda y \left(1 - \frac{\lambda}{3}\right)^2 \right)^{\frac{1}{4}} \right\} d\lambda \\ = 2 \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (1, 1), (\xi, 1); \\ (1, 4), (\xi + \eta, 1); \end{matrix} \frac{4y}{9} \right] \end{aligned} \quad (3.7)$$

Corollary 8 If we put $c = 1$ in equation (2.2) and use equ (1.7), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho,\sigma} \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (\rho, \sigma k), (\eta, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.8)$$

Corollary 9 If we put $c = \sigma = 1$ in equation (2.2) and use equ (1.6), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{k,r,s}^{\rho} \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (\rho, k), (\eta, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.9)$$

Corollary 10 If we put $c = \sigma = k = 1$ in equation (2.2) and use equ (1.5), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{r,s}^{\rho} \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (\rho, 1), (\eta, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.10)$$

Corollary 11 If we put $c = \sigma = k = \rho = 1$ in equation (2.2) and use equ (1.4), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_{r,s} \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (1, 1), (\eta, 1); \\ (s, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.11)$$

Corollary 12 If we put $c = \sigma = k = \rho = s = 1$ in equation (2.2) and use equ (1.1), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} E_r \left[y(1-\lambda)^2 \left(1 - \frac{\lambda}{4}\right) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (1, 1), (\eta, 1); \\ (1, r), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.12)$$

Corollary 13 If we put $c = \sigma = k = \rho = s = 1$ and $r = 2$ in equation (2.2) and use equ (1.2), we get

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \cosh \left[\sqrt{y \left(1 - \frac{\lambda}{4}\right)} (1-\lambda) \right] d\lambda \\ = \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta)_2 \psi_2 \left[\begin{matrix} (1, 1), (\eta, 1); \\ (1, 2), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.13)$$

Corollary 14 *If we put $c = \sigma = k = \rho = s = 1$ and $r = 4$ in equation (2.2) and use equ (1.3), we get*

$$\begin{aligned} \int_0^1 \lambda^{\xi-1} (1-\lambda)^{2\eta-1} \left(1 - \frac{\lambda}{3}\right)^{2\xi-1} \left(1 - \frac{\lambda}{4}\right)^{\eta-1} \left\{ \cos [y(1-\lambda)^2 (1 - \frac{\lambda}{4})]^{\frac{1}{4}} + \cosh [y(1-\lambda)^2 (1 - \frac{\lambda}{4})]^{\frac{1}{4}} \right\} d\lambda \\ = 2 \left(\frac{2}{3}\right)^{2\xi} B(\xi, \eta) {}_2\psi_2 \left[\begin{matrix} (1, 1), (\eta, 1); \\ (1, 4), (\xi + \eta, 1); \end{matrix} y \right] \end{aligned} \quad (3.14)$$

4. Conclusion

By involving extended k -generalized Mittag-Leffler function, two unified integrals established in the form of generalized Wright hyper-geometric function. Also discussed certain cases which are converted into the elementary Mittag-Leffler function.

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